

CMP Revision

- Binding
 - Lattice + Basis
 - Reciprocal Lattice & Diffraction
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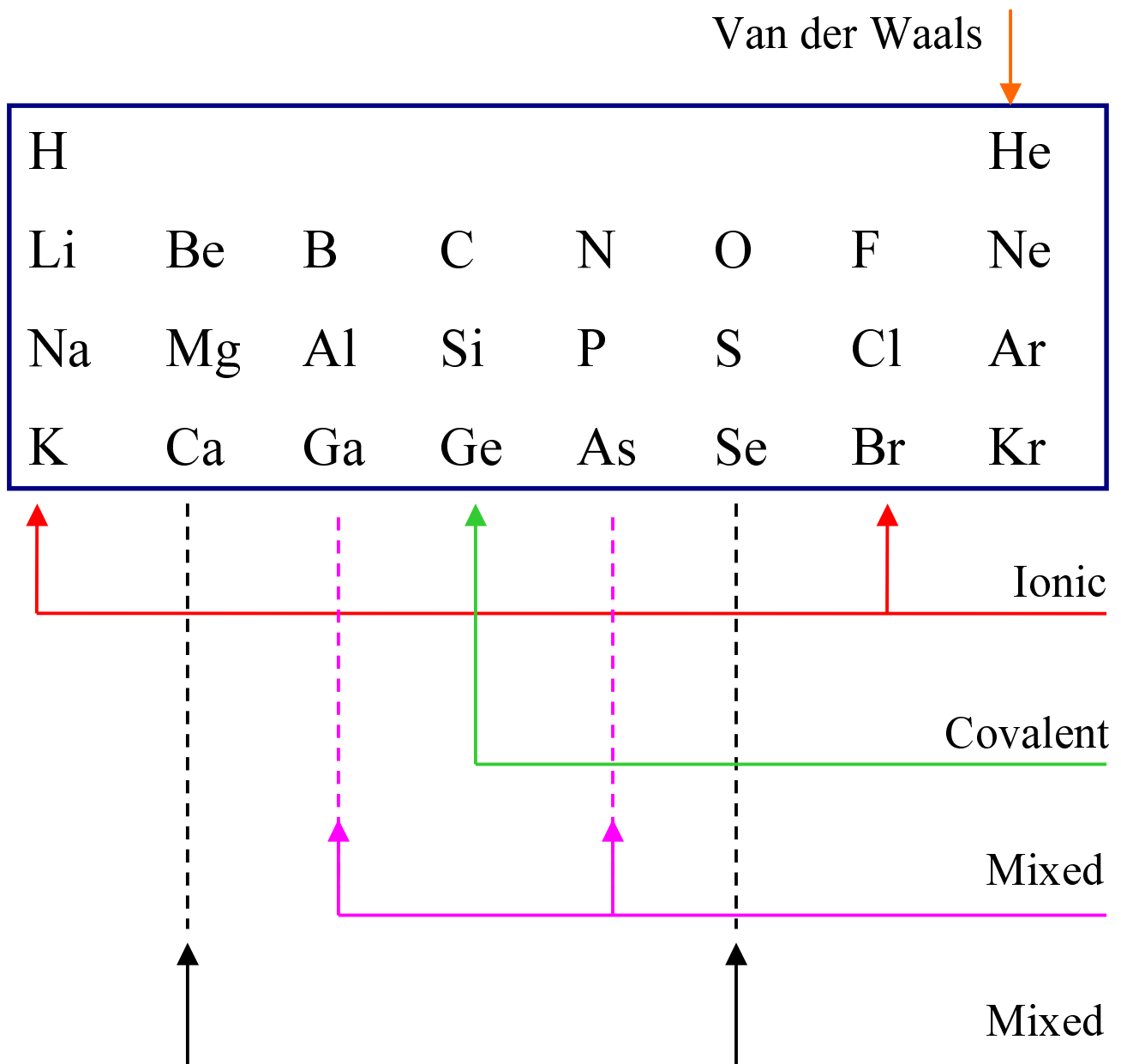
Excitations – \mathbf{k} in a periodic structure

- Phonons
 - Lattice Specific Heat
 - Thermal Conductivity
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Superconductivity

- Type I
- Behaviour in Field
- Paired electrons
- Flux quantization

Types of Binding

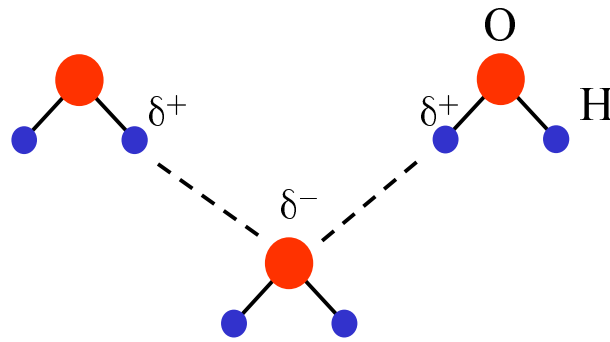


Ionic (e.g. NaCl) → Na^+Cl^- **Strong**

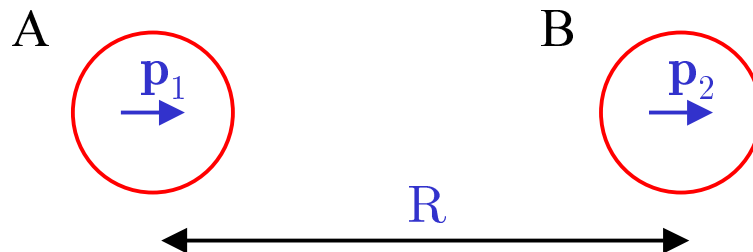
Covalent (e.g. C) → **Strong**

In covalent bonds, spins are paired anti-symmetrically, whilst orbital Ψ s are symmetric.

Hydrogen Bond - Ice



Van der Waals forces



Instantaneous dipole moment \mathbf{p}_1 on A produces an electric field

$|\mathbf{E}| \propto \frac{|\mathbf{p}_1|}{\epsilon_0 R^3}$ at B. This induces \mathbf{p}_2 on B given by:

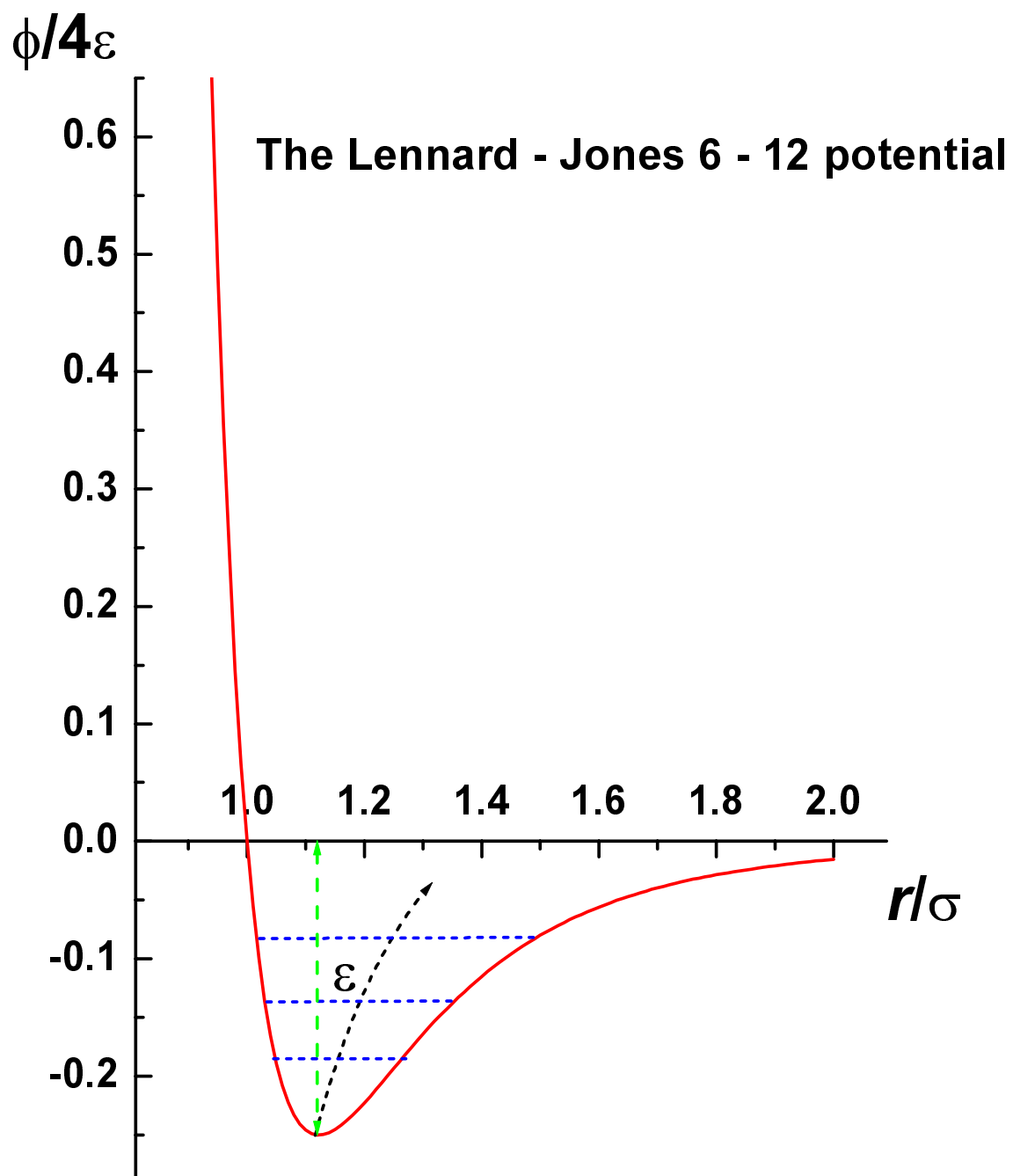
$|\mathbf{p}_2| = \alpha |\mathbf{E}|$ (where α is the polarizability)

The energy of interaction of two dipoles is:

$$U \propto \frac{p_1 p_2}{R^3} \propto \frac{\alpha p_1^2}{R^6}$$

The core electrons repel each other strongly when the wavefunctions begin to overlap. Experimentally it is found that the repulsion $\sim 1/R^{12}$. The full potential is then:

$$\phi(r) = 4\epsilon \left[\left(\frac{\sigma}{r} \right)^{12} - \left(\frac{\sigma}{r} \right)^6 \right]$$



For ionic crystals we have:

$$\phi(r) = -\frac{\alpha e^2}{4\pi\epsilon_0 r} + \frac{C}{r^m}$$

Where α is known as the Madelung constant.

What about Bulk Modulus? Using the bulk modulus and r_0 we can calculate m .

$$B = -V \left(\frac{\partial P}{\partial V} \right)_T$$

and at $T = 0$

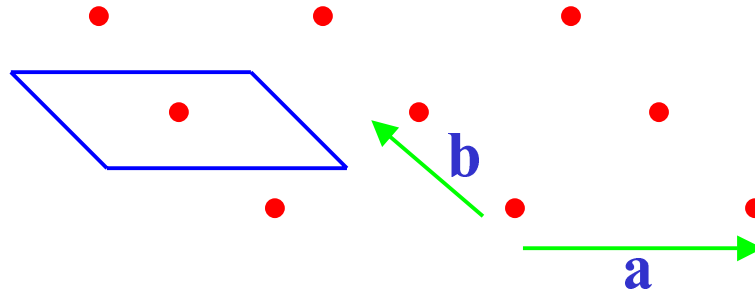
$$P = -\frac{dU}{dV}$$

now for $u = U / N$ and $v = V / N$

$$B = v \frac{\partial}{\partial v} \left(\frac{\partial u}{\partial v} \right) \text{ and } v = f(r^3)$$

See page 402 Ashcroft and Mermin.

What makes a crystal? – Long range order

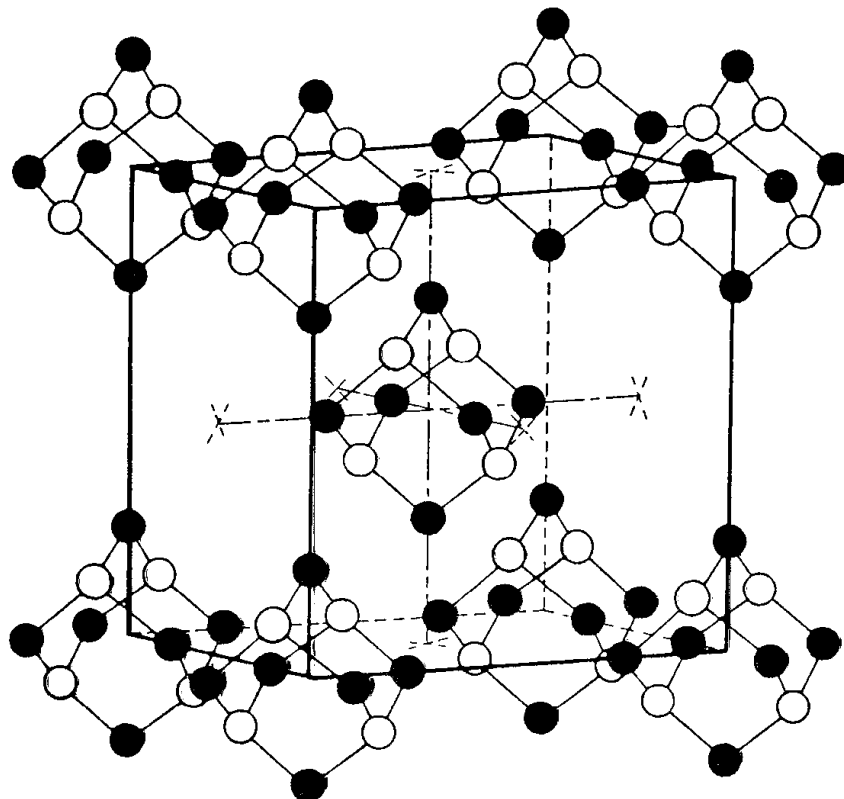


Translational symmetry

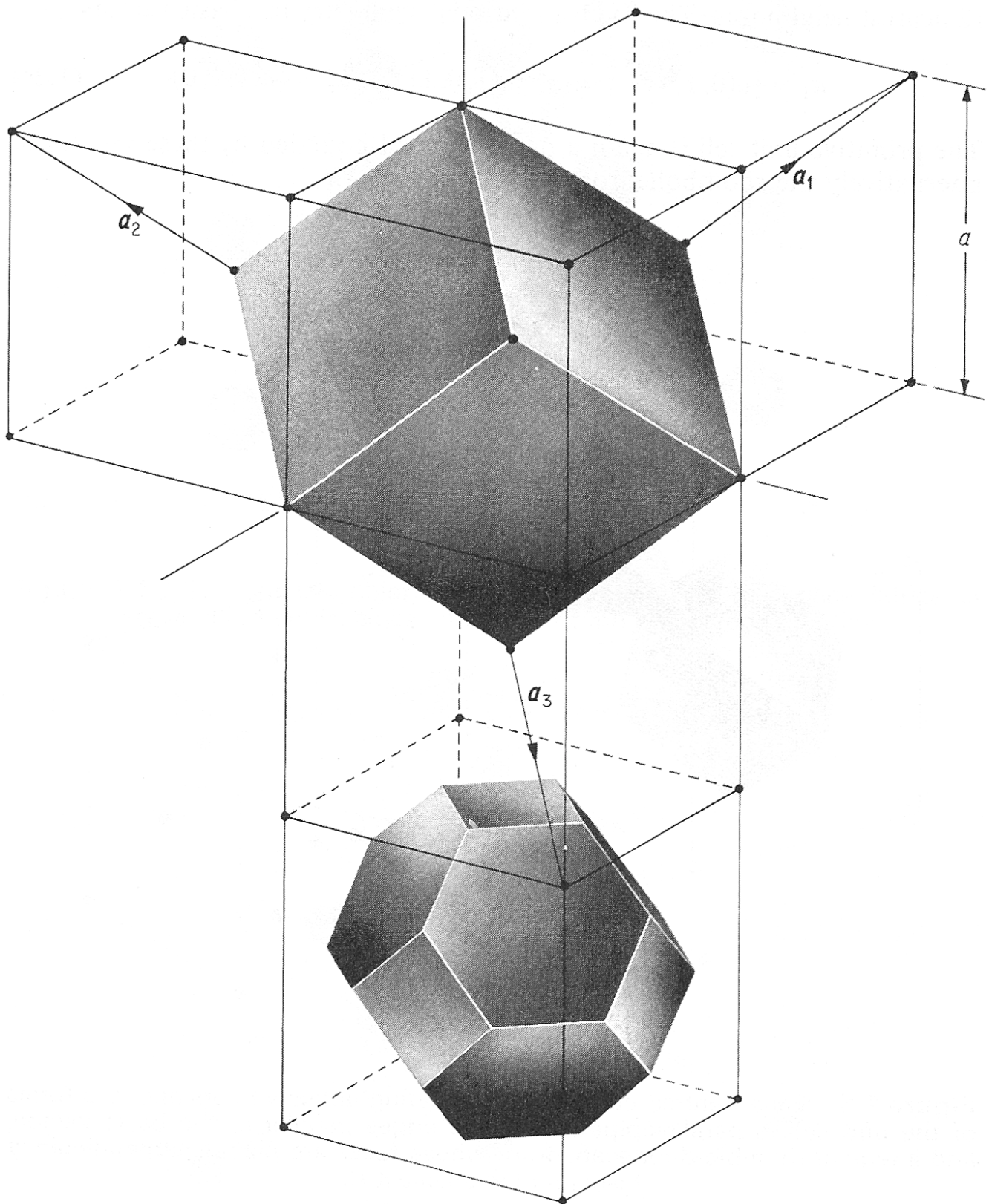
$$P(\mathbf{r}) \equiv P(\mathbf{r} + \mathbf{T})$$

$$\text{where } \mathbf{T} = n_1 \mathbf{a}_1 + n_2 \mathbf{a}_2 + n_3 \mathbf{a}_3$$

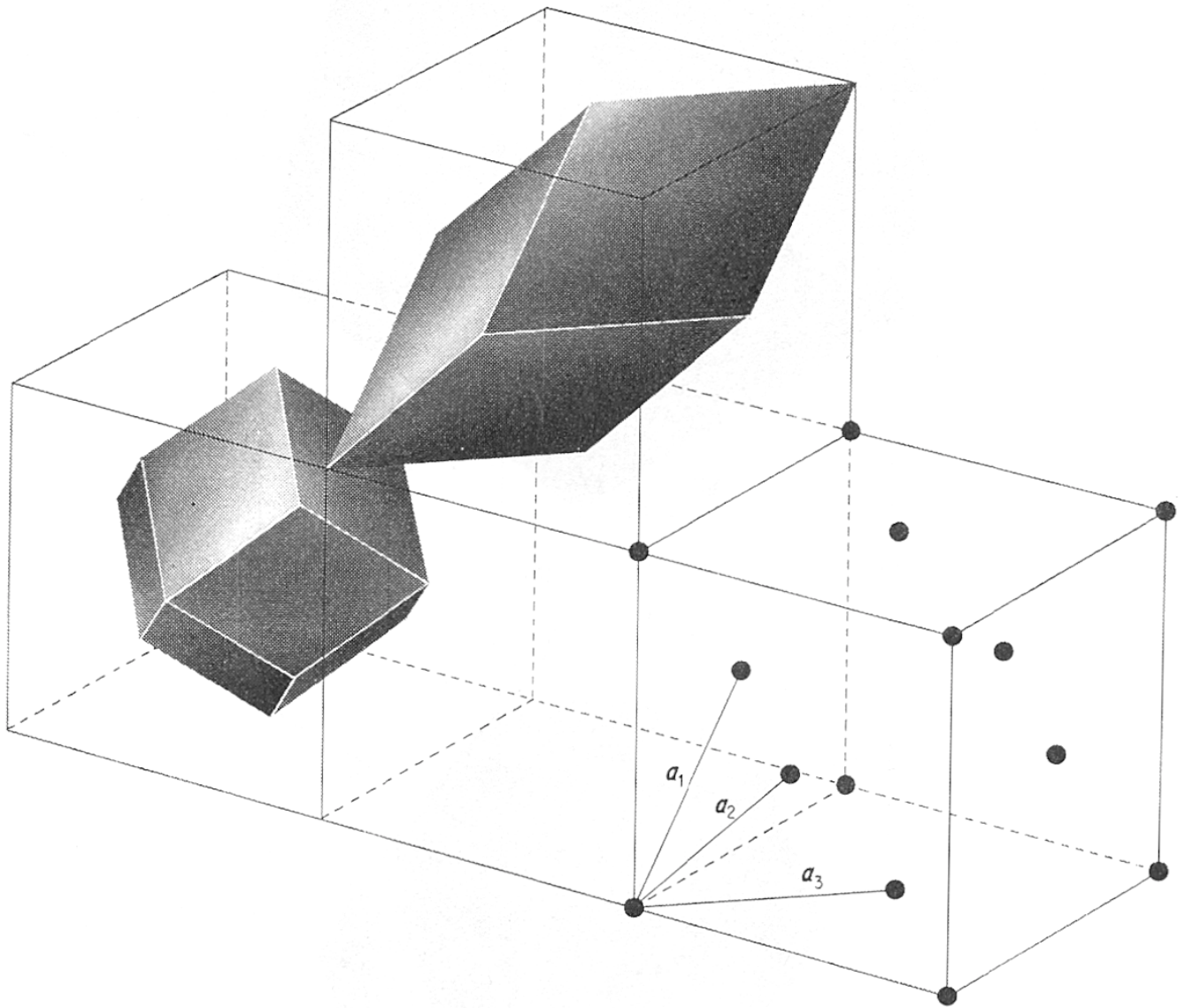
This defines a **LATTICE** of points in space, for full structure we need a **BASIS** of atoms at each point.



The crystal structure of hexamethylene tetramine $\text{C}_6\text{H}_{12}\text{N}_4$ is based on a body-centred cubic lattice. Each carbon (black in diagram) is bonded to two hydrogens, which are not shown.



Body-centred cubic lattice, showing basic vectors, and two forms of the unit cell, a parallelepiped with the basic vectors as edges and a truncated octahedron where the faces are the planes perpendicularly bisecting the smallest R_i



Face-centred cubic lattice, showing lattice vectors and two forms of the unit cell.

Reciprocal Lattice

Properties of a crystal such as electron density and electrostatic potential are periodic and satisfy:

$$F(\mathbf{r} + \mathbf{T}) = F(\mathbf{r})$$

By analogy with 1D periodic functions we expand in a Fourier series:

$$F(\mathbf{r}) = \sum_{\mathbf{G}} A(\mathbf{G}) e^{i\mathbf{G} \cdot \mathbf{r}}$$

Where \mathbf{G} is such that $e^{i\mathbf{G} \cdot \mathbf{r}}$ is also periodic in \mathbf{T} ,

$$\therefore e^{i\mathbf{G} \cdot \mathbf{T}} = 1$$

$$\text{or } \mathbf{G} \cdot \mathbf{T} = 2\pi \times \text{integer}$$

Vectors \mathbf{G} define the **Reciprocal Lattice**.

$$\text{Let } \mathbf{G} = m_1 \mathbf{g}_1 + m_2 \mathbf{g}_2 + m_3 \mathbf{g}_3$$

$$\text{Then } \mathbf{G}_i \cdot \mathbf{T}_i = 2\pi(n_1 m_1 + n_2 m_2 + n_3 m_3)$$

$$\text{in 3D, } \mathbf{g}_1 = \frac{2\pi \mathbf{a}_2 \wedge \mathbf{a}_3}{\mathbf{a}_1 \cdot (\mathbf{a}_2 \wedge \mathbf{a}_3)}$$

$$\text{or } \mathbf{a}_i \cdot \mathbf{g}_j = 2\pi \delta_{ij}$$

Some Results To Be Aware Of

1. The R.L.V. $\mathbf{G} = h\mathbf{A} + k\mathbf{B} + l\mathbf{C}$ is perpendicular to a plane (hkl) in the crystal lattice.
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2. The distance between two adjacent parallel planes of the lattice (planes through lattice points) is :

$$d(hkl) = \frac{2\pi}{|\mathbf{G}|}$$

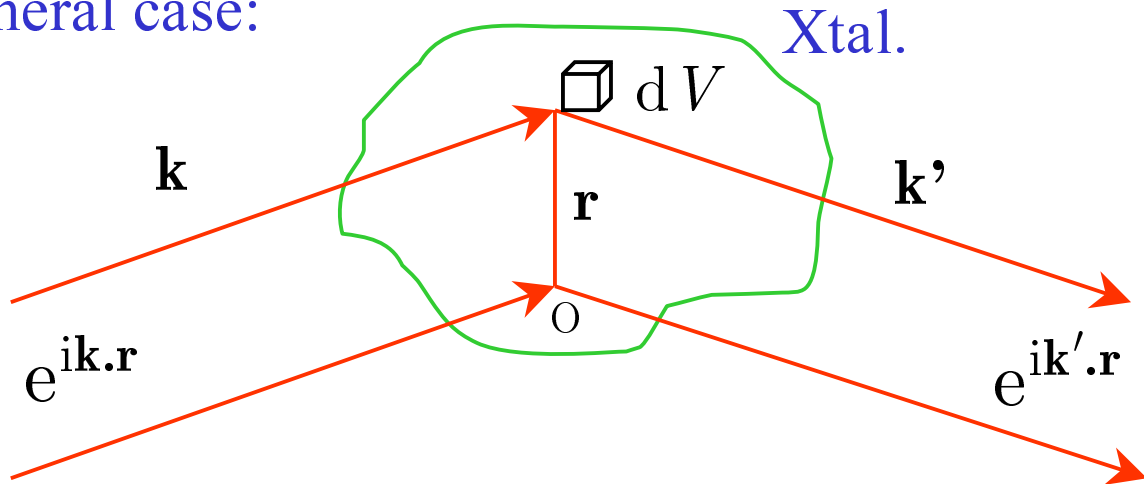
3. For a cubic lattice:

$$d^2 = \frac{a^2}{(h^2 + k^2 + l^2)}$$

Where a is the cubic side.

Diffraction

General case:



Difference in phase: $e^{i(\mathbf{k}-\mathbf{k}') \cdot \mathbf{r}}$

Total scattered amplitude:

$$\mathbf{A} = \int dV \, n(\mathbf{r}) e^{-i(\Delta \mathbf{k} \cdot \mathbf{r})}$$

Taking the Fourier Transform gives:

$$\begin{aligned} \mathbf{A} &= \sum \int dV \, n_{\mathbf{G}} \exp[i(\mathbf{G} - \Delta \mathbf{k}) \cdot \mathbf{r}] \\ &= 0 \quad \text{unless} \quad \mathbf{G} = \Delta \mathbf{k} \quad \text{or} \end{aligned}$$

$$2\mathbf{k} \cdot \mathbf{G} + \mathbf{G}^2 = 0$$

or

$$2\left(\frac{2\pi}{\lambda}\right) \sin \theta = |\mathbf{G}| = \frac{2\pi}{d}$$

$$\therefore 2d \sin \theta = n\lambda$$

As each cell is the same,

$$\mathbf{A}_{\mathbf{G}} = N \int_{\text{cell}} dV \, n(\mathbf{r}) e^{-i\mathbf{G} \cdot \mathbf{r}} = N \mathbf{S}_{\mathbf{G}}$$

$\mathbf{S}_{\mathbf{G}}$ is the **Structure Factor**. Referring electron concentrations to one corner of each cell,

$$n(\mathbf{r}) = \sum_{j=1}^s n_j (\mathbf{r} - \mathbf{r}_j) = \sum_{j=1}^s n_j (\bar{\rho})$$

over s atoms of the basis. Therefore:

$$\mathbf{S}_{\mathbf{G}} = \sum_j e^{-i\mathbf{G} \cdot \mathbf{r}_j} \int dV \, n_j(\bar{\rho}) e^{-i\mathbf{G} \cdot \bar{\rho}}$$

$$\& \mathbf{S}_{\mathbf{G}} = \sum_j f_j e^{-i\mathbf{G} \cdot \mathbf{r}_j}$$

where f_j is the **Form Factor**

$$\text{if} \quad \mathbf{r}_j = x_j \mathbf{a}_1 + y_j \mathbf{a}_2 + z_j \mathbf{a}_3$$

$$\& \quad \mathbf{G} = h\mathbf{b}_1 + k\mathbf{b}_2 + l\mathbf{b}_3$$

$$S(hkl) = \sum_j f_j \exp[-i2\pi(x_j h + y_j k + z_j l)]$$

& the scattered intensity $\propto S^* S$

Results to remember

BCC

$$S = 0 \quad (h + k + l) \text{ odd}$$

$$S = 2f \quad (h + k + l) \text{ even}$$

FCC

$$S = 4f \quad (h, k, l) \text{ all odd or all even}$$

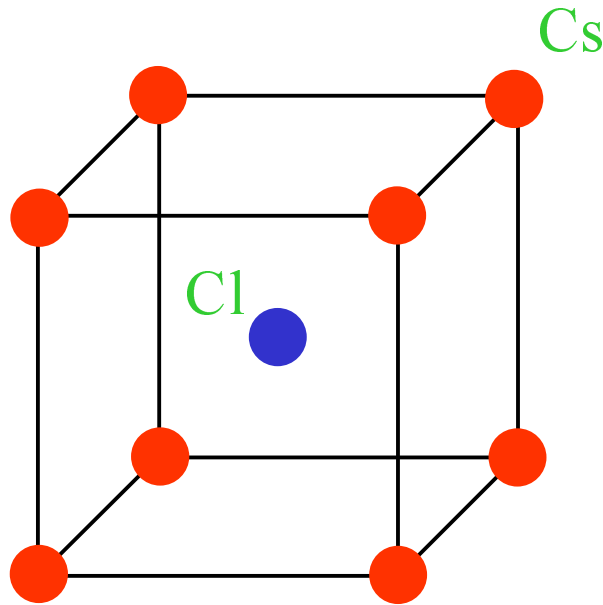
$$S = 0 \quad \text{otherwise}$$

N.B. For FINALS QUESTIONS

Simple Cubic	Body C. Cubic	Face C. Cubic
100	_____	_____
110	110	_____
111	_____	111
200	200	200
etc.	etc.	etc.

What if the Atoms are Different?

Example CsCl



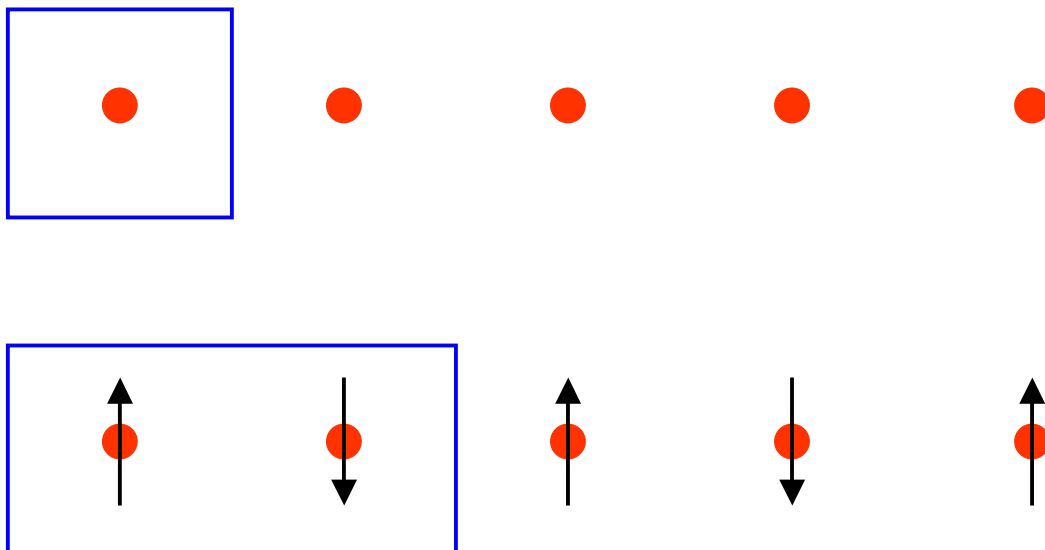
For (100), $S = f_{\text{Cs}} - f_{\text{Cl}}$

As f_{Cs} and f_{Cl} are not equal,
(100) is present.

Probes for Diffraction

- **X-rays:** scatter off electrons
Probability of scattering $\sim Z^3$ Therefore
useless for small Z atoms (e.g. hydrogen,
lithium)
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- **Neutrons:** have mass
 - (a) nuclear interactions via strong force
 - (b) magnetic interactions via
magnetic moment



- **Electrons:** good for surfaces
LEED, RHEED, BLEED.

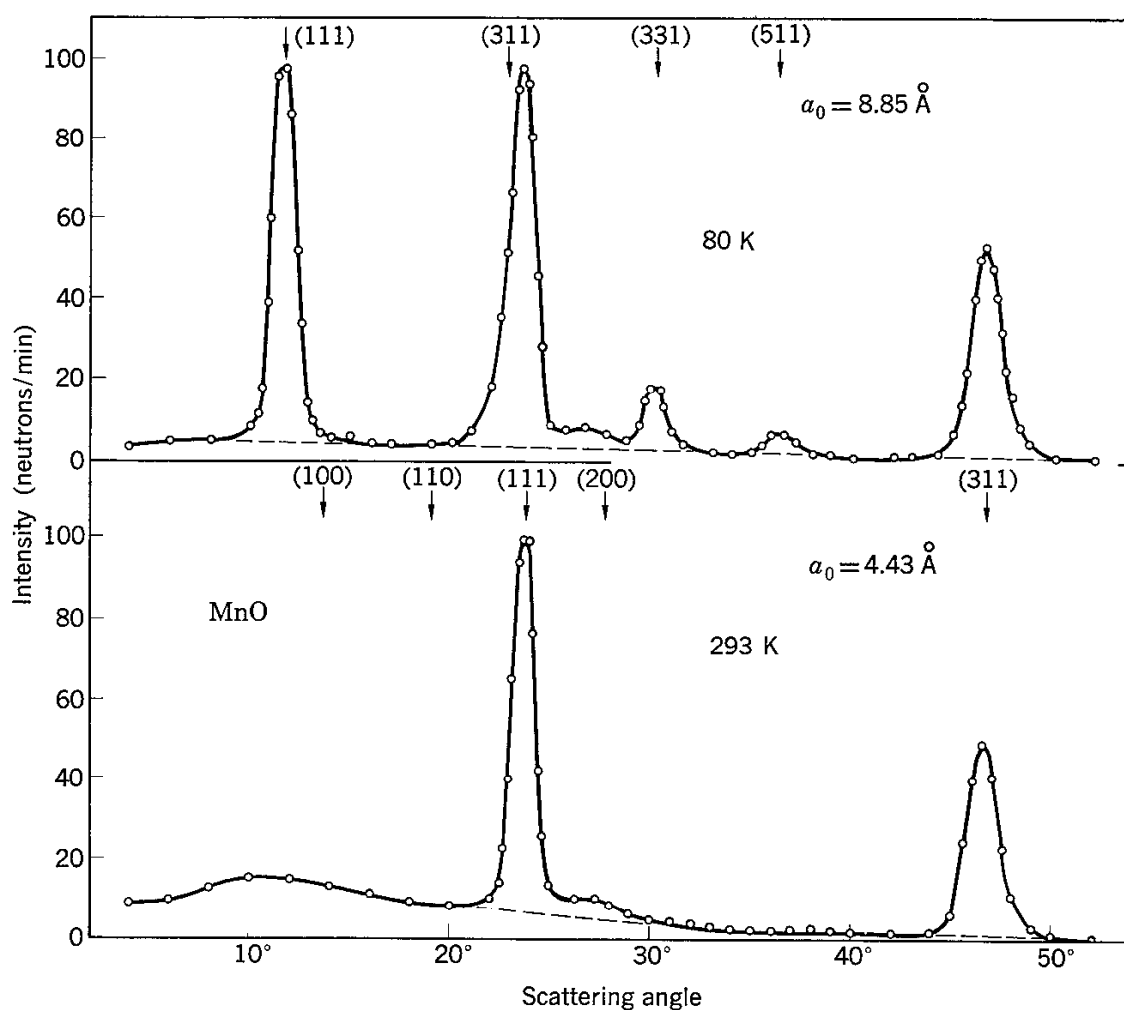
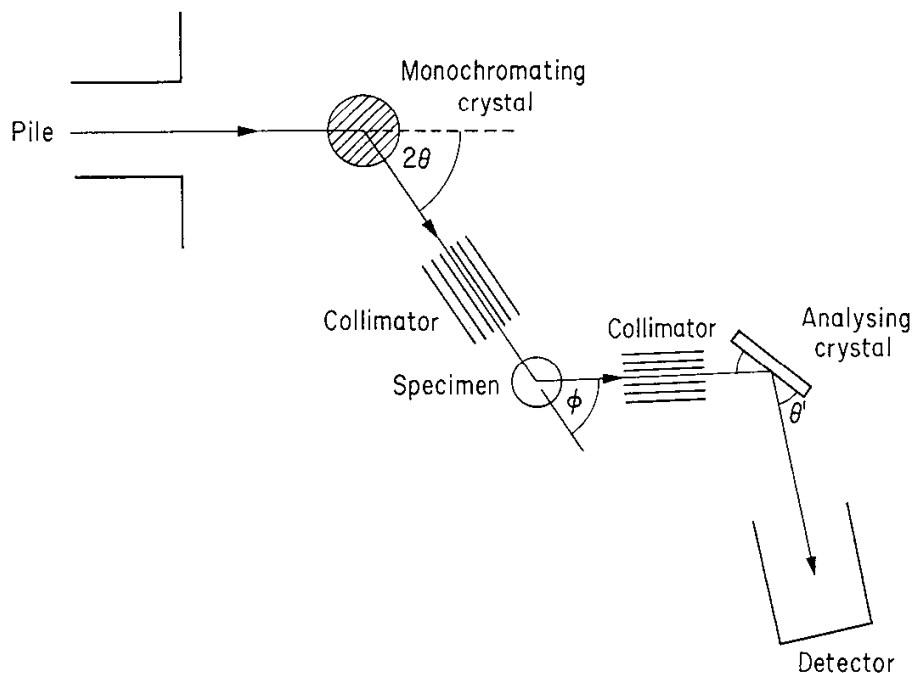


Figure 20 Neutron diffraction patterns for MnO below and above the spin-ordering temperature of 120 K, after C. G. Shull, W. A. Strauser, and E. O. Wollan. Phys. Rev. **83**, 333 (1951). The reflection indices are based on an 8.85 Å cell at 80 K and on a 4.43 Å cell at 293 K. At the higher temperature the Mn^{2+} ions are still magnetic, but they are no longer ordered.