

D'Alembert's solution

$$\frac{\partial^2 y}{\partial x^2} = \frac{1}{c^2} \frac{\partial^2 y}{\partial t^2} \quad \Rightarrow \quad y(x, t) = f(x - ct) + g(x + ct)$$

f and g are determined by the initial conditions:

Suppose at time $t = 0$, the wave has an initial displacement $U(x)$ and an initial velocity $V(x)$

$$y(x, 0) = f(x) + g(x) = U(x)$$

$$\frac{\partial y(x, 0)}{\partial t} = -cf'(x) + cg'(x) = V(x) \quad \Rightarrow \quad f(x) - g(x) = -\frac{1}{c} \int_b^x V(x') dx'$$

$$f(x) = \frac{1}{2}U(x) - \frac{1}{2c} \int_b^x V(x') dx'$$

$$g(x) = \frac{1}{2}U(x) + \frac{1}{2c} \int_b^x V(x') dx'$$

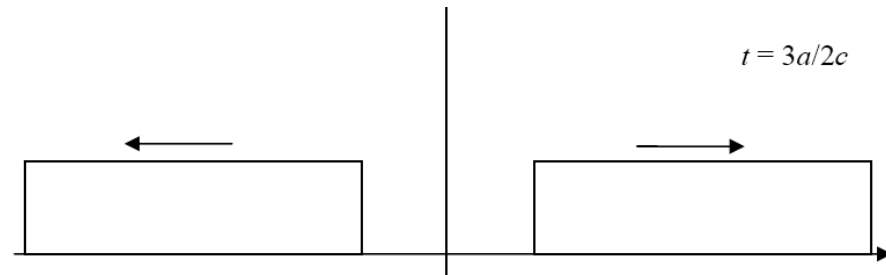
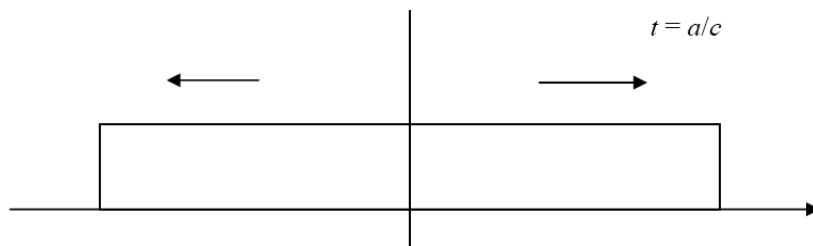
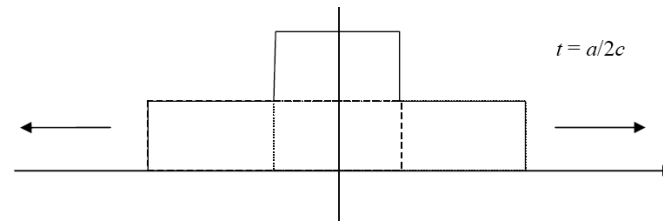
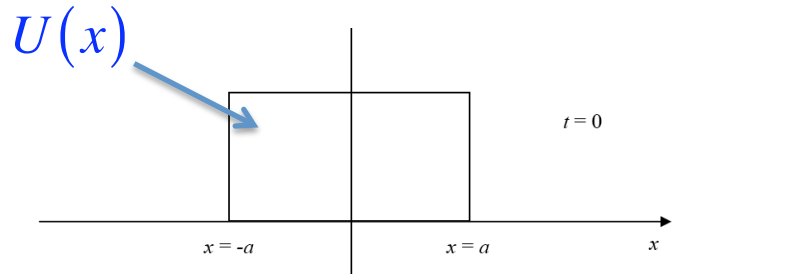
$$y(x, t) = \frac{1}{2} [U(x - ct) + U(x + ct)] + \frac{1}{2c} \left[\int_b^{x+ct} V(x) dx - \int_b^{x-ct} V(x) dx \right] = \frac{1}{2} [U(x - ct) + U(x + ct)] + \frac{1}{2c} \int_{x-ct}^{x+ct} V(x) dx$$

Ex. Wave with initial rectangular displacement released from rest,

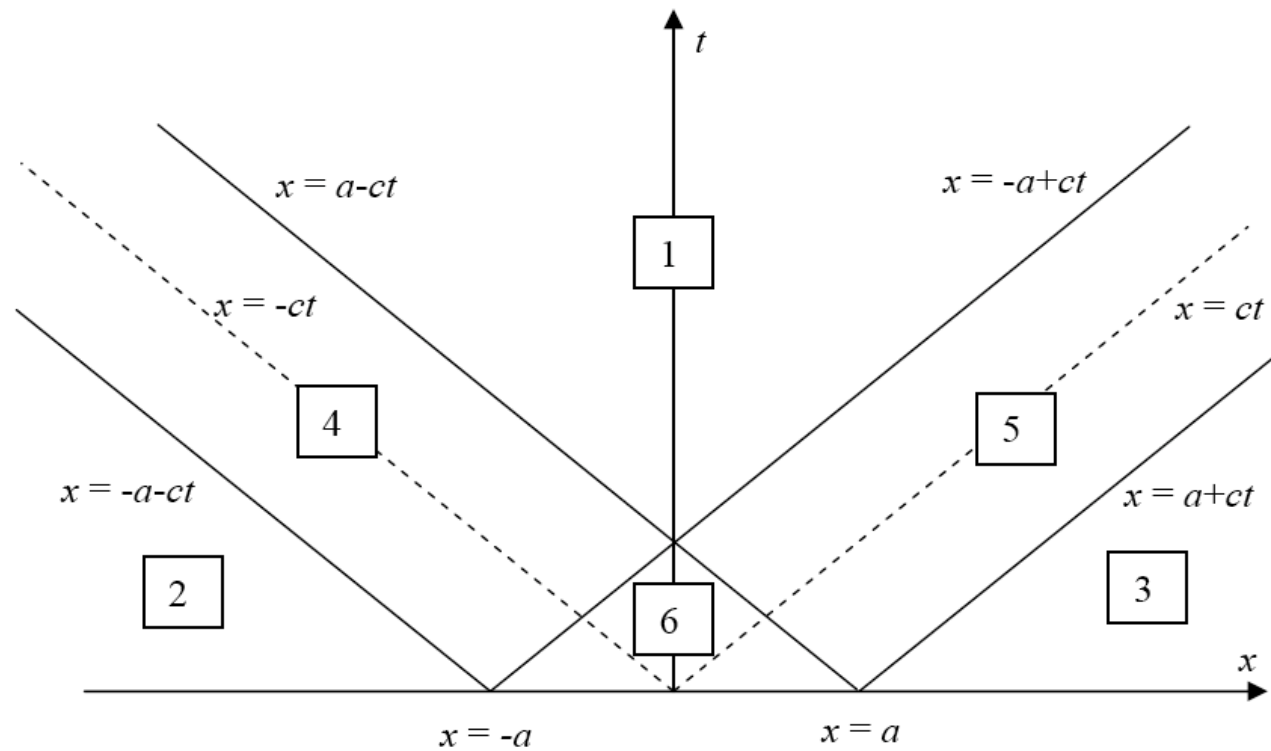
$$V(x) = 0$$

$$y(x,t) = \frac{1}{2} [U(x-ct) + U(x+ct)]$$

$$\left(y(x,t) = \frac{1}{2} [U(x-ct) + U(x+ct)] + \frac{1}{2c} \int_{x-ct}^{x+ct} V(x) dx \right)$$



This figure can also be represented on an (x,t) domain. Let $y(x,t)$ be pointing out of the paper.



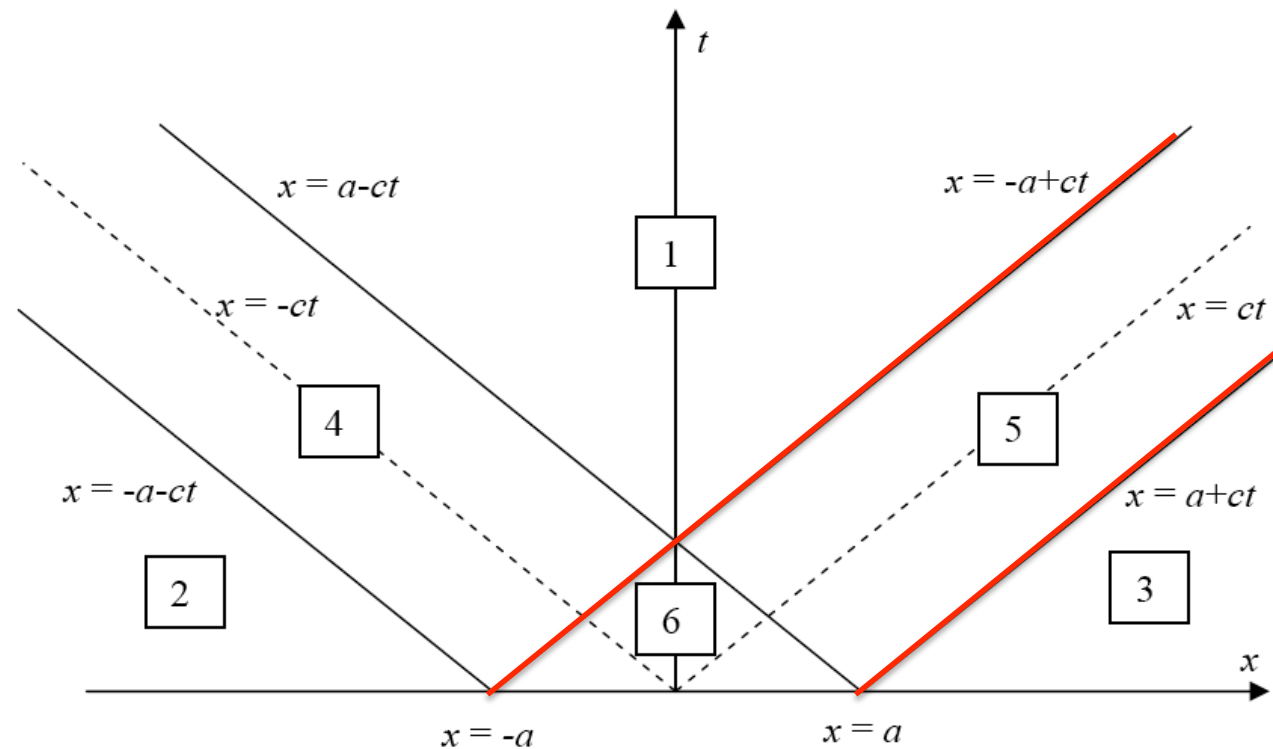
In regions 1, 2, and 3, $y(x,t) = 0$ for all x, t

$$\text{In region 5, } y(x,t) = \frac{1}{2}u(x-ct) \quad -a \leq x-ct \leq a$$

$$\text{In region 4, } y(x,t) = \frac{1}{2}u(x+ct) \quad -a \leq x+ct \leq a$$

$$\text{In region 6, } y(x,t) = \frac{1}{2}[u(x-ct) + u(x+ct)] \quad \begin{array}{l} x-ct > -a \\ x+ct < a \end{array}$$

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In region 4, $y(x,t) = \frac{1}{2}u(x+ct)$ $-a \leq x+ct \leq a$

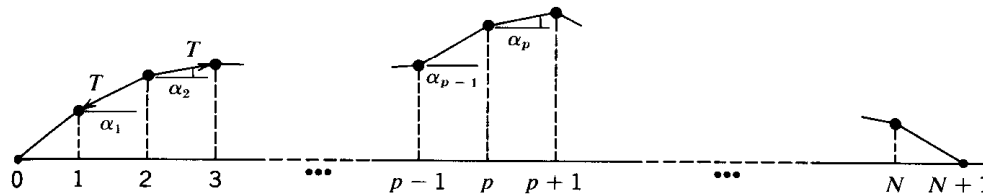
In region 6, $y(x,t) = \frac{1}{2}[u(x-ct) + u(x+ct)]$ $x-ct > -a$
 $x+ct < a$

$$y(x,t) = f(x - ct) + g(x + ct)$$

What is the form of $f(x)$, $g(x)$?

N coupled oscillators

Consider the transverse oscillations of N particles of mass m spaced equally along a flexible, elastic, massless string, which is under tension T .



Assume the particles are displaced by small distances y_i :

Consider the p^{th} particle

$$F_y = -T \sin \alpha_{p-1} + T \sin \alpha_p \approx -\frac{T}{l}(y_p - y_{p-1}) + \frac{T}{l}(y_{p+1} - y_p)$$

$$\ddot{y}_p + 2\omega_0^2 y_p - \omega_0^2 (y_{p+1} - y_{p-1}) = 0, \quad \omega_0^2 = T / ml$$

General solution :

$$y_p = \sum_{n=1}^N \sin\left(\frac{pn\pi}{N+1}\right) \left(D_n \cos(\omega_n t + \delta_n) \right)$$

Superposition of N normal modes

$$\omega_n = 2\omega_0 \sin\left[\frac{n\pi}{2(N+1)}\right]$$

$$y(x,t) = f(x - ct) + g(x + ct)$$

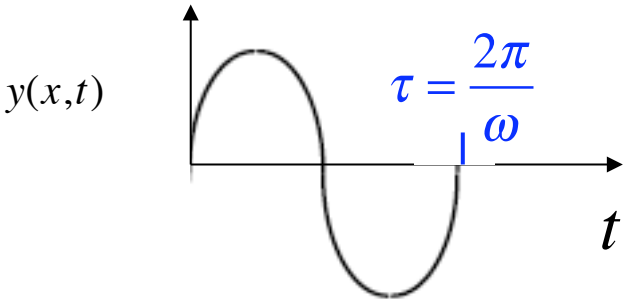
If time dependence is $\cos(\omega t)$ the full (x,t) dependence is given by

$$y(x,t) = A \cos(kx + \omega t) + B \cos(kx - \omega t)$$

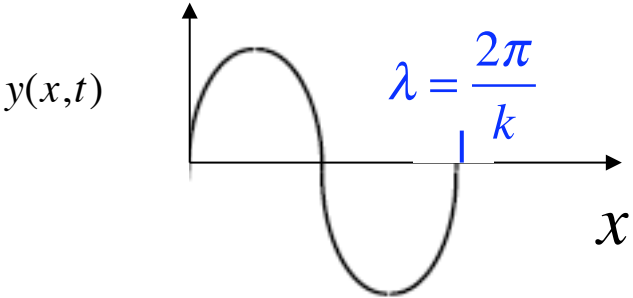
constant k

- Speed of wave $c = \frac{\omega}{k}$

- Frequency $f = \frac{1}{\tau} = \frac{\omega}{2\pi}$



- Wavelength $\lambda = \frac{2\pi}{k}$



k is "wavenumber"

We can write the equation of a travelling wave in a number of analogous forms:

	Velocity	Wavelength	Period	Angular frequency
$A \sin(kx - \omega t)$	ω / k	$2\pi / k$	$2\pi / \omega$	ω
$A \sin k(x - vt)$	v	$2\pi / k$	$2\pi / vk$	vk
$A \sin \left[2\pi \left(\frac{x}{\lambda} - \frac{t}{\tau} \right) \right]$	λ / τ	λ	τ	$2\pi / \tau$
$A \sin \left[2\pi (x - vt) / \lambda \right]$	v	λ	λ / v	$2\pi v / \lambda$

N.B. Can include phase most easily by putting

$$y(x, t) = \text{Re} \left[A \exp \left[i(kx - \omega t) \right] \right]$$

where A is complex.

N.B.2 Sometimes more convenient to switch x and t , i.e.

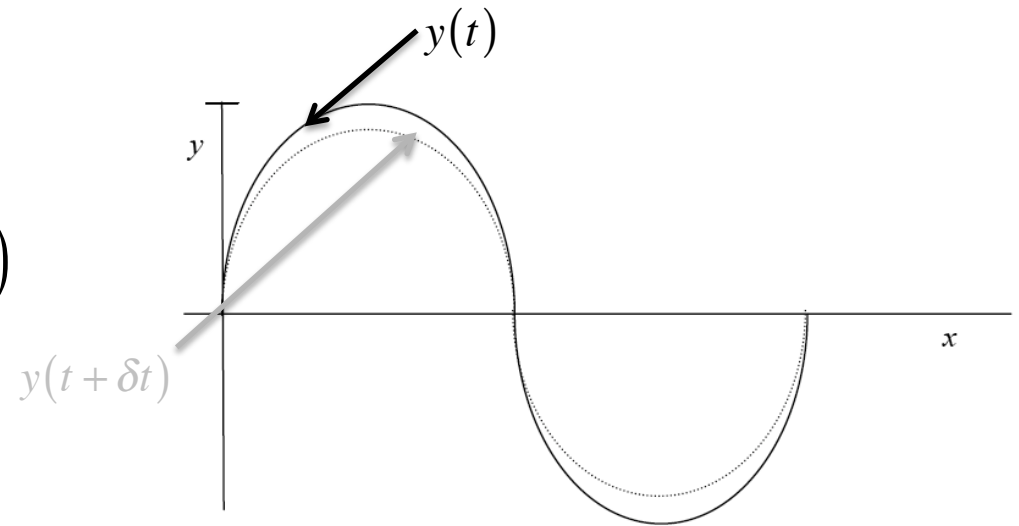
$$y(x, t) = A \sin(\omega t - kx)$$

This is still a travelling wave moving to the right.

For non-sinusoidal wave moving to right with speed v , can always write as $f(x - vt)$.

Stationary waves

$$\begin{aligned}y &= A \sin(kx - \omega t) + A \sin(kx + \omega t) \\ &= 2A \sin kx \cos \omega t\end{aligned}$$



“stationary wave”

More generally

$$\begin{aligned}y &= A \sin(kx - \omega t + 2\delta_1) + A \sin(kx + \omega t + 2\delta_2) \\ &= 2A \sin(kx + \delta_1 + \delta_2) \cos(\omega t + \delta_1 - \delta_2)\end{aligned}$$