

# Advanced Quantum Mechanics

## Scattering Theory and Relativistic Quantum Mechanics

### Problem Set 1: Scattering Theory

- 1 Find the Born approximation to the differential cross section as a function of the angle of scattering for a central potential  $V(r) = -V_0$  for  $r \leq r_0$ ,  $V(r) = 0$  for  $r > r_0$ .
- 2 The potential for ‘hard sphere’ scattering is  $V(r) = \infty$  for  $r < r_0$  and  $V(r) = 0$  for  $r > r_0$ .
  - (i) Show that the s-wave phase shift  $\delta_0(k)$  is given by  $\delta_0(k) = -kr_0$ , and calculate the total  $\ell = 0$  cross section.
  - (ii) The ‘reduced’ radial wave function  $u_\ell(r)$  is defined via  $\psi(r, \theta, \phi) = R_\ell(r)Y_{\ell,m}(\theta, \phi)$ , and  $R_\ell(r) = u_\ell(r)/r$ . Find the equation satisfied by  $u_\ell$ , and show that the radial wave equation for  $r > r_0$  and  $\ell = 1$  is satisfied by the wave function

$$u_1(r) = C \left[ \frac{\sin kr}{kr} - \cos kr + a \left( \frac{\cos kr}{kr} + \sin kr \right) \right]$$

where  $C$  and  $a$  are constants. Show that the definition of the p-wave phase shift  $\delta_1(k)$  implies that  $\tan \delta_1(k) = a$ , and determine  $a$  from the condition satisfied by  $u_1(r)$  at  $r = r_0$ .

Show that, as  $k$  approaches zero,  $\delta_1(k)$  behaves like  $(kr_0)^3$ .

- 3 The scattering amplitude  $f(\theta, k)$  has the partial wave expansion

$$f(\theta, k) = \frac{1}{k} \sum_{\ell} e^{i\delta_\ell} \sin \delta_\ell (2\ell + 1) P_\ell(\cos \theta)$$

where the symbols have their usual meaning.

(i) Show

(a)  $\sigma_{tot} = \sum_{\ell} \sigma_{\ell}$  where  $\sigma_{\ell} = \frac{4\pi}{k^2} \sin^2 \delta_{\ell} (2\ell + 1)$

(b)  $\sigma_{tot} = \frac{4\pi}{k} \text{Im} f(0, k)$ .

(ii) In the ‘Fermi theory’ of lepton-lepton scattering the s-wave amplitude  $a_0 \equiv e^{i\delta_0} \sin \delta_0 / k$  for the scattering process

$$\bar{\nu}_\mu + \mu^- \rightarrow \bar{\nu}_e + e^-$$

is given approximately by  $a_0^F \approx G_F k$  where  $k$  is the centre of mass momentum and  $G_F$  is Fermi’s constant, with the value (in units  $\hbar = c = 1$ )

$$G_F \approx \frac{10^{-5}}{M_p^2}$$

where  $M_p$  is the proton mass. Assuming that  $k$  in part (i) can also be taken to be the CM momentum (in units  $c = \hbar = 1$ ) estimate the CM energy (neglecting lepton masses) at which unitarity is violated in the theory. What actually happens? (Note:  $\hbar \approx 200(\text{MeV}/c)\text{fm}$ ).

- 4  $L_x$  is a linear ordinary differential operator and  $\psi(x)$  satisfies the equation  $L_x\psi(x) = g(x)$  where  $g(x)$  is a known function. The functions  $\phi_n(x)$  form a complete orthonormal set of eigenfunctions of  $L_x$ , with eigenvalues  $\lambda_n$ :

$$L_x\phi_n(x) = \lambda_n\phi_n(x).$$

Expanding  $\psi(x)$  as a linear combination of the  $\phi_n(x)$  via

$$\psi(x) = \sum_n a_n\phi_n(x),$$

show that

$$a_m = \frac{1}{\lambda_m} \int \phi_m^*(y)g(y)dy.$$

Deduce that the Green function for the equation  $L_x\psi = g$  must be

$$G(x, y) = \sum_n \frac{\phi_n(x)\phi_n^*(y)}{\lambda_n} = \sum_n \frac{\langle x|n\rangle \langle n|y\rangle}{\lambda_n}.$$

Verify that this is correct by applying the operator  $L_x$  to  $G(x, y)$ .

Use this approach to find the coordinate space Green function for the time independent Schrodinger equation.

- 5 (i) In first order time-dependent perturbation theory, the transition probability per unit time for a particle in the momentum eigenstate  $|\mathbf{k}_i\rangle$  to scatter, via the time-independent potential  $V(r)$ , into the momentum eigenstate  $|\mathbf{k}_f\rangle$  is

$$\int 2\pi |\langle \mathbf{k}_f | V | \mathbf{k}_i \rangle|^2 \delta(E_i - E_f) \frac{d^3\mathbf{k}_f}{(2\pi)^3}$$

where units  $\hbar = 1$  are used (for a concise derivation see P79, Quarks and Leptons, Halzen & Martin, Wiley 1984). Writing

$$d^3\mathbf{k}_f = |\mathbf{k}_f|^2 d|\mathbf{k}_f| d\Omega \quad \text{and} \quad E_f = \mathbf{k}_f^2/2m$$

show that  $d^3\mathbf{k}_f = m |\mathbf{k}_f| dE_f d\Omega$ , and hence show that the differential cross section is

$$\frac{d\sigma}{d\Omega} = \frac{m^2}{4\pi^2} |\langle f | V | i \rangle|^2.$$

- (ii) Show that

$$\int e^{i\mathbf{q}\cdot\mathbf{r}} \frac{e^{-r/a}}{r} d^3\mathbf{r} = \frac{4\pi}{\mathbf{q}^2 + a^{-2}}.$$

By considering the limit  $a \rightarrow \infty$ , derive the Rutherford scattering cross section formula

$$\left(\frac{d\sigma}{d\Omega}\right)_R = \frac{\hbar^2 c^2 \alpha^2}{16E^2 \sin^4 \theta/2}$$

using Born Approximation for the Coulomb potential for two charges, one  $+e$  and one  $-e$ ; here  $E = \frac{1}{2}mv^2$  where  $v$  is the speed of the incoming particle and  $m$  is its mass; also,  $\alpha$  is the fine structure constant and  $\theta$  is the scattering angle (ie the angle between the incident wave vector and the final wave vector).