

Exponential $h(x)$

$$Lf = a_2 \frac{d^2 f}{dx^2} + a_1 \frac{df}{dx} + a_0 f = h(x) = He^{\gamma x}$$

$$\text{CF} = A_{1,2} e^{\alpha_{1,2} x}.$$

$$\equiv a_2 \left(\frac{d}{dx} - \alpha_1 \right) \left(\frac{d}{dx} - \alpha_2 \right) f = He^{\gamma x}.$$

$$\text{PI: } f = \frac{He^{\gamma x}}{a_2(\gamma - \alpha_1)(\gamma - \alpha_2)}$$

$$\alpha_1 \neq \alpha_2 \neq \gamma \neq \alpha_1$$

$$f = \frac{Hxe^{\alpha_2 x}}{a_2(\alpha_2 - \alpha_1)}$$

$$\gamma = \alpha_2 \neq \alpha_1$$

$$\text{CF} = Ae^{\alpha x} + Bxe^{\alpha x}$$

$$\alpha_1 = \alpha_2 = \alpha = \gamma$$

$$\text{PI: } \text{Try } f = Px^2 e^{\alpha x}$$

$$a_2 \left(\frac{d}{dx} - \alpha_1 \right) \left(\frac{d}{dx} - \alpha_2 \right) f = H e^{\gamma x}.$$

Try $f = P x^2 e^{\alpha x}$

$$a_2 \left(\frac{d}{dx} - \alpha \right)^2 P x^2 e^{\alpha x} = a_2 \left(\frac{d}{dx} - \alpha \right) 2 P x e^{\alpha x} = 2 a_2 P e^{\alpha x} = H e^{\alpha x}$$

$$\Rightarrow \boxed{P = \frac{H}{2 a_2}}$$

Exponential $h(x)$

$$Lf = a_2 \frac{d^2 f}{dx^2} + a_1 \frac{df}{dx} + a_0 f = h(x).$$

$$CF = A_{1,2} e^{\alpha_{1,2} x}.$$

PI :

$$f = \frac{He^{\gamma x}}{a_2(\gamma - \alpha_1)(\gamma - \alpha_2)}$$

$$\alpha_1 \neq \alpha_2 \neq \gamma \neq \alpha_1$$

$$f = \frac{Hxe^{\alpha_2 x}}{a_2(\alpha_2 - \alpha_1)}$$

$$\gamma = \alpha_2 \neq \alpha_1$$

$$f = \frac{H}{2a_2} x^2 e^{\alpha x}$$

$$\alpha_1 = \alpha_2 = \alpha = \gamma$$

Ex 4

$$f'' + 3f' + 2f = e^{-x}, \quad f(0) = 1, \quad f'(0) = 1$$

The **CF** is $f_0 = Ae^{-2x} + Be^{-x}$

PI ... try $f_1 = Pxe^{-x}$

$$\begin{aligned} e^{-x} &= \left(\frac{d}{dx} + 2\right)\left(\frac{d}{dx} + 1\right)Pxe^{-x} = \left(\frac{d}{dx} + 2\right)Pe^{-x} \\ &= Pe^{-x}. \end{aligned}$$

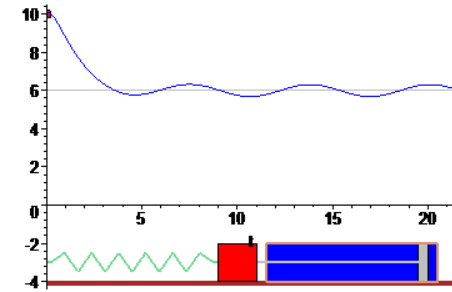
$$P = 1 \quad \text{and} \quad f_1 = xe^{-x}$$

Full solution : $f = f_0 + f_1 = xe^{-x} + Ae^{-2x} + Be^{-x}$

Initial conditions \Rightarrow $f = xe^{-x} - e^{-2x} + 2e^{-x}$

Oscillators

$$m\ddot{x} = - \underbrace{m\omega_0^2 x}_{\text{spring}} - \underbrace{m\gamma\dot{x}}_{\text{friction}} + \underbrace{mF \cos \omega t}_{\text{forcing}}$$



The associated complex equation is

$$\ddot{z} + \gamma\dot{z} + \omega_0^2 z = F e^{i\omega t}.$$

Transients :

CF - Auxiliary equation : $z = e^{\alpha t}$.

$$\alpha^2 + \gamma\alpha + \omega_0^2 = 0 \Rightarrow \alpha = -\frac{1}{2}\gamma \pm i\sqrt{\omega_0^2 - \frac{1}{4}\gamma^2}$$

$$= -\frac{1}{2}\gamma \pm i\omega_\gamma \quad \text{where} \quad \omega_\gamma \equiv \omega_0 \sqrt{1 - \frac{1}{4}\gamma^2/\omega_0^2}.$$

$$\omega_0^2 - \frac{1}{4}\gamma^2 > 0$$

Complementary function

$$x = e^{-\gamma t/2} [A \cos(\omega_\gamma t) + B \sin(\omega_\gamma t)] = e^{-\gamma t/2} N \cos(\omega_\gamma t + \psi)$$

Constant phase "shift"

Since $\gamma > 0$, the CF $\rightarrow 0$ as $t \rightarrow \infty$ CF describes "transients"

Steady state solutions

$$\ddot{z} + \gamma\dot{z} + \omega_0^2 z = F e^{i\omega t}.$$

No damping
exponential

Particular integral

$$x = \Re e\left(\frac{F e^{i\omega t}}{\omega_0^2 - \omega^2 + i\omega\gamma}\right).$$

i.e. the Particular integral describes the steady state solution after the transients have died away.

Since the denominator = $\sqrt{(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2} e^{i\phi}$ where $\phi \equiv \arctan\left(\frac{\omega\gamma}{\omega_0^2 - \omega^2}\right)$ the

particular integral can be written as

$$x = \frac{F \Re e(e^{i(\omega t - \phi)})}{\sqrt{(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2}} = \frac{F \cos(\omega t - \phi)}{\sqrt{(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2}}$$

For $\phi > 0$, x achieves the same phase as F at t greater by $\left(\frac{\phi}{\omega}\right)$

- ϕ is called the "phase lag" of the response.

$$x = \frac{F \Re(e^{i(\omega t - \phi)})}{\sqrt{(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2}} = \frac{F \cos(\omega t - \phi)}{\sqrt{(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2}}$$

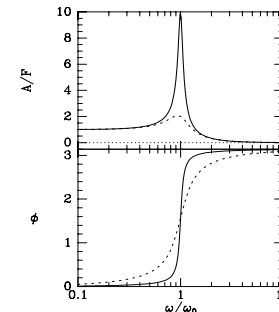
The amplitude of the response is $A = \frac{F}{\sqrt{(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2}}$,

This has a maximum when $0 = \frac{dA^{-2}}{d\omega} \propto -4(\omega_0^2 - \omega^2)\omega + 2\omega\gamma^2 \Rightarrow \omega^2 = \omega_0^2 - \frac{1}{2}\gamma^2$.

$\omega_R \equiv \sqrt{\omega_0^2 - \gamma^2/2}$ is called the “resonant” frequency

The frictional coefficient causes the resonant frequency to be less than the normal frequency

$$\phi \equiv \arctan\left(\frac{\omega\gamma}{\omega_0^2 - \omega^2}\right)$$



Oscillators

$$m\ddot{x} = \underbrace{-m\omega_0^2 x}_{\text{spring}} - \underbrace{m\gamma\dot{x}}_{\text{friction}} + \underbrace{mF \cos \omega t}_{\text{forcing}}$$

$$\mathbf{F} = mF \cos \omega t, \quad x = \frac{F \Re e(e^{i(\omega t - \phi)})}{\sqrt{(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2}} = \frac{F \cos(\omega t - \phi)}{\sqrt{(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2}}$$

Power Input

(steady state) $P = \frac{\partial W}{\partial t} = \mathbf{F}\dot{x}$ $W = \int_{x(t_0)}^{x(t)} F dx'$

$$\begin{aligned} P = \mathbf{F}\dot{x} &= mF \cos \omega t \times \frac{-F\omega \sin(\omega t - \phi)}{\sqrt{(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2}} \\ &= \frac{\omega m F^2}{\sqrt{(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2}} [-\cos(\omega t) \sin(\omega t - \phi)] \\ &= -\frac{\frac{1}{2} \omega m F^2}{\sqrt{(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2}} [\sin(2\omega t - \phi) + \sin(-\phi)]. \end{aligned}$$

Average over a period

$$\overline{P} = \frac{\frac{1}{2} \omega m F^2 \sin \phi}{\sqrt{(\omega_0^2 - \omega^2)^2 + \omega^2 \gamma^2}}.$$

Energy dissipated

$$m\ddot{x} = \underbrace{-m\omega_0^2 x}_{\text{spring}} - \underbrace{m\gamma\dot{x}}_{\text{friction}} + \underbrace{mF \cos \omega t}_{\text{forcing}}.$$

$$\overline{D} = m\gamma \overline{\dot{x}\dot{x}} = \frac{m\gamma\omega^2 F^2 \frac{1}{2}}{(\omega_0^2 - \omega^2)^2 + \omega^2\gamma^2}.$$

$$x = \frac{F \cos(\omega t - \phi)}{\sqrt{(\omega_0^2 - \omega^2)^2 + \omega^2\gamma^2}}$$

$$\overline{\sin^2(\omega t - \phi)} = \frac{1}{2}$$

$$\overline{D} = \overline{P} = \frac{\frac{1}{2}\omega m F^2 \sin \phi}{\sqrt{(\omega_0^2 - \omega^2)^2 + \omega^2\gamma^2}} \quad \text{since} \quad \sin \phi = \gamma\omega / \sqrt{(\omega_0^2 - \omega^2)^2 + \omega^2\gamma^2}$$

$$\left(\phi \equiv \arctan\left(\frac{\omega\gamma}{\omega_0^2 - \omega^2}\right) \right)$$

Quality Factor

$$m\ddot{x} = \underbrace{-m\omega_0^2 x}_{\text{spring}} - \underbrace{m\gamma\dot{x}}_{\text{friction}} + \underbrace{mF \cos \omega t}_{\text{forcing}}$$

Energy content of transient motion that the CF describes

$$E = \frac{1}{2}(m\dot{x}^2 + m\omega_0^2 x^2)$$

$$x = e^{-\gamma t/2} A \cos(\omega_\gamma t + \psi)$$

$$= \frac{1}{2} m A^2 e^{-\gamma t} \left[\frac{1}{4} \gamma^2 \cos^2 \eta + \omega_\gamma \gamma \cos \eta \sin \eta + \omega_\gamma^2 \sin^2 \eta + \omega_0^2 \cos^2 \eta \right] \quad (\eta \equiv \omega_\gamma t + \psi)$$

$$E \simeq \frac{1}{2} m (\omega_0 A)^2 e^{-\gamma t} \quad \text{small } \frac{\gamma}{\omega_0} \quad (\omega_\gamma \equiv \omega_0 \sqrt{1 - \frac{1}{4} \gamma^2 / \omega_0^2} \simeq \omega_0)$$

Quality factor

$$Q \equiv \frac{E(t)}{E(t - \pi/\omega_0) - E(t + \pi/\omega_0)} \simeq \frac{1}{e^{\pi\gamma/\omega_0} - e^{-\pi\gamma/\omega_0}} = \frac{1}{2} \operatorname{csc} h(\pi\gamma/\omega_0)$$

$$\simeq \frac{\omega_0}{2\pi\gamma} \quad (\text{for small } \gamma/\omega_0).$$

Q is the inverse of the fraction of the oscillator's energy that is dissipated in one period
 - approximately the number of oscillations before the energy decays by factor e