

The Standard Model (& Beyond)

G. Ross, May 2009

The Standard Model

$$L_{SM} = L_{YM} + L_{WD} + L_{Yu} + L_H$$

Strong, weak and EM
interactions

A renormalisable, spontaneously broken, local gauge quantum field theory

The Standard Model

$SU(3) \times SU(2) \times U(1)$

$$L_{SM} = L_{YM} + L_{WD} + L_{Yu} + L_H$$

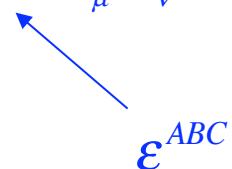
$$L_{YM} = L_{QCD} + L_{I_W} + L_Y$$

$$= -\frac{1}{4g_3^2} \sum_{A=1}^8 G_{\mu\nu}^A G^{\mu\nu A} - \frac{1}{4g_2^2} \sum_{a=1}^3 F_{\mu\nu}^a F^{\mu\nu a} - \frac{1}{4g_1^2} B_{\mu\nu} B^{\mu\nu}$$

$$G_{\mu\nu}^A = \partial_\mu A_\nu^A - \partial_\nu A_\mu^A - f^{ABC} A_\mu^B A_\nu^C, \quad A, B, C = 1, \dots, 8.$$

$$[T^A, T^B] = if^{ABC} T^C$$

$$F_{\mu\nu}^A = \partial_\mu W_\nu^A - \partial_\nu W_\mu^A - f^{ABC} W_\mu^B W_\nu^C, \quad A, B, C = 1, \dots, 3.$$



$$\epsilon^{ABC}$$

The Standard Model

$$L_{SM} = L_{YM} + L_{WD} + L_{Yu} + L_H$$

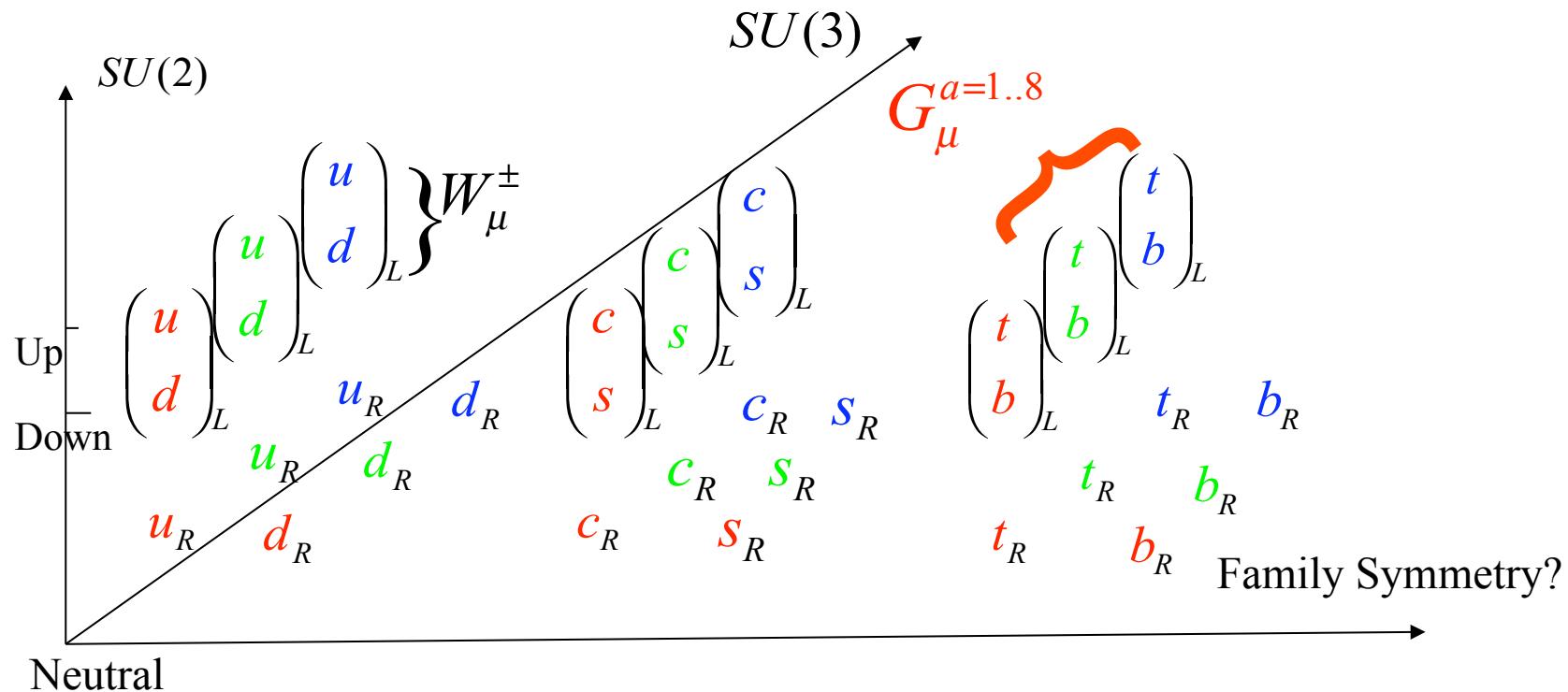
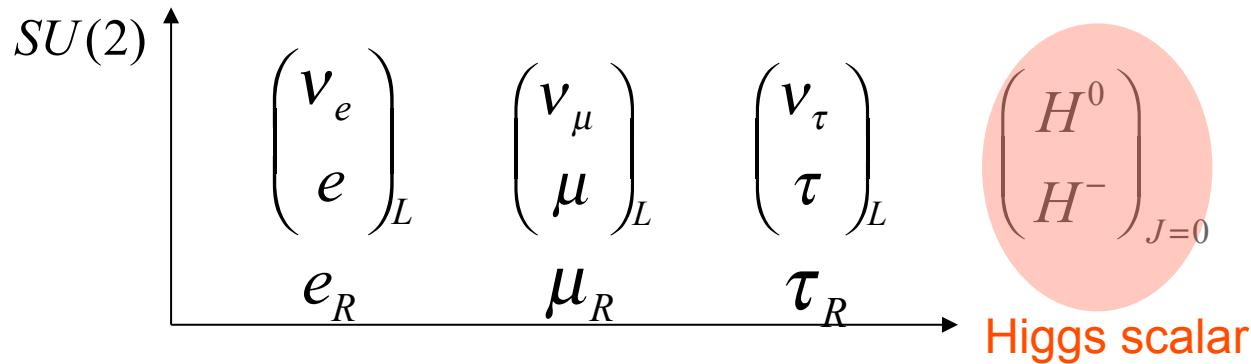
$$L_{YM} = L_{QCD} + L_{I_W} + L_Y$$

$$= -\frac{1}{4g_3^2} \sum_{A=1}^8 G_{\mu\nu}^A G^{\mu\nu A} + \frac{\theta}{32\pi^2} \epsilon^{\mu\nu\lambda\sigma} \sum_{A=1}^8 G_{\mu\nu}^A G_{\lambda\sigma}^A - \frac{1}{4g_2^2} \sum_{a=1}^3 F_{\mu\nu}^a F^{\mu\nu a} - \frac{1}{4g_1^2} B_{\mu\nu} B^{\mu\nu}$$

$\theta < 10^{-7}$ strong CP problem

Multiplet structure -'chiral'

$$SU(3) \otimes SU(2) \otimes U(1)$$



$$L_{WD} = \sum_{i=1}^3 \left(L_i^\dagger \sigma^\mu D_\mu L_i + \bar{e}_i^\dagger \sigma^\mu D_\mu \bar{e}_i + Q_i^\dagger \sigma^\mu D_\mu Q_i + \bar{u}_i^\dagger \sigma^\mu D_\mu \bar{u}_i + \bar{d}_i^\dagger \sigma^\mu D_\mu \bar{d}_i \right)$$

$$D_\mu L_i = (\partial_\mu + iW_\mu + \frac{i}{2}y_1 B_\mu) L_i$$

$$y_1 = -1$$

$$D_\mu \bar{e}_i = (\partial_\mu + \frac{i}{2}y_2 B_\mu) \bar{e}_i$$

$$y_2 = +2$$

$$D_\mu Q_i = (\partial_\mu + iA_\mu + iW_\mu + \frac{i}{2}y_3 B_\mu) Q_i$$

$$y_3 = +\frac{1}{3}$$

?

$$D_\mu \bar{u}_i = (\partial_\mu - iA_\mu + \frac{i}{2}y_4 B_\mu) \bar{u}_i$$

$$y_4 = -\frac{4}{3}$$

$$D_\mu \bar{d}_i = (\partial_\mu - iA_\mu + \frac{i}{2}y_5 B_\mu) \bar{d}_i$$

$$y_5 = +\frac{2}{3}$$

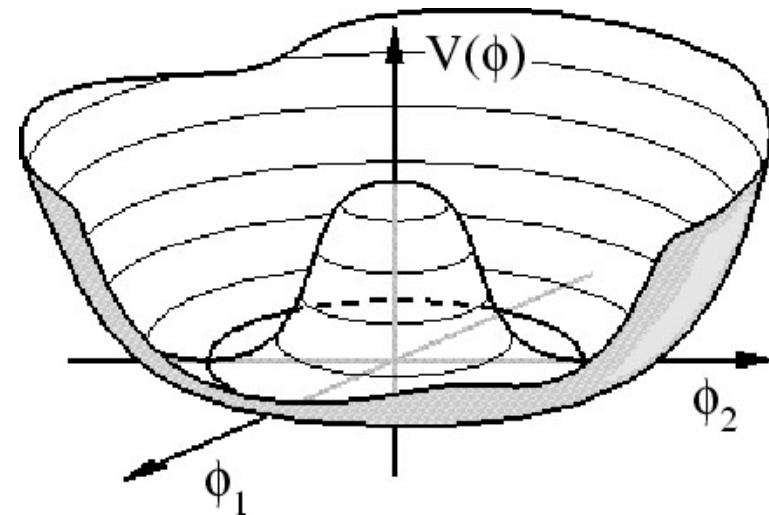
$$L_{Yu} = i \hat{L}_i \bar{e}_j H^* Y_{ij}^l + i \hat{Q}_i \bar{d}_j H^* Y_{ij}^d + i \hat{Q}_i \bar{u}_j \tau_2 H Y_{ij}^u + c.c.$$

Spontaneously broken

$$L_H = \left(D_\mu H \right)^\dagger \left(D^\mu H \right) - V(H)$$

$$V = -m^2 |H|^2 + \lambda |H|^4$$

$$H = e^{i\xi\tau} \begin{pmatrix} 0 \\ \frac{v+h}{\sqrt{2}} \end{pmatrix}$$

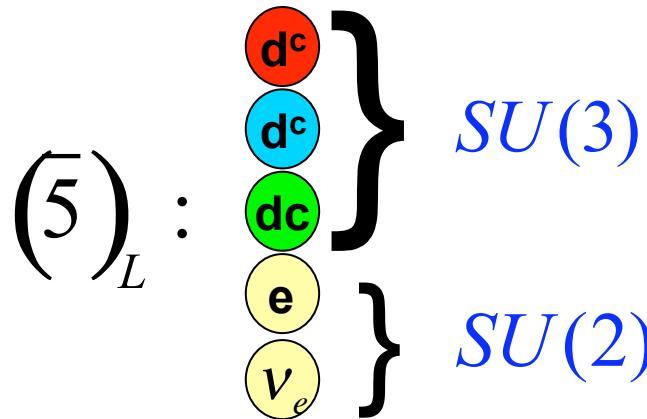


$$L_H = \frac{1}{2} \partial_\mu h \partial^\mu h + \frac{1}{2} M_Z^2 Z_\mu Z^\mu + \frac{1}{2} M_W^2 W_\mu^+ W^{-\mu} + \frac{h}{v} \left(2 + \frac{h}{v} \right) \left(\frac{1}{2} M_Z^2 Z_\mu Z^\mu + \frac{1}{2} M_W^2 W_\mu^+ W^{-\mu} \right)$$

The Standard Model (& Beyond?)

Grand Unification

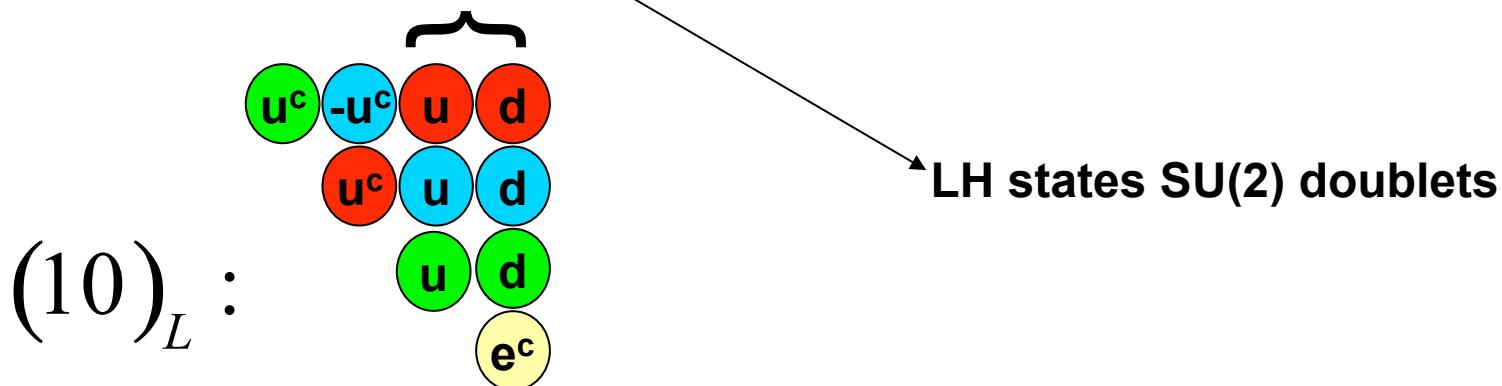
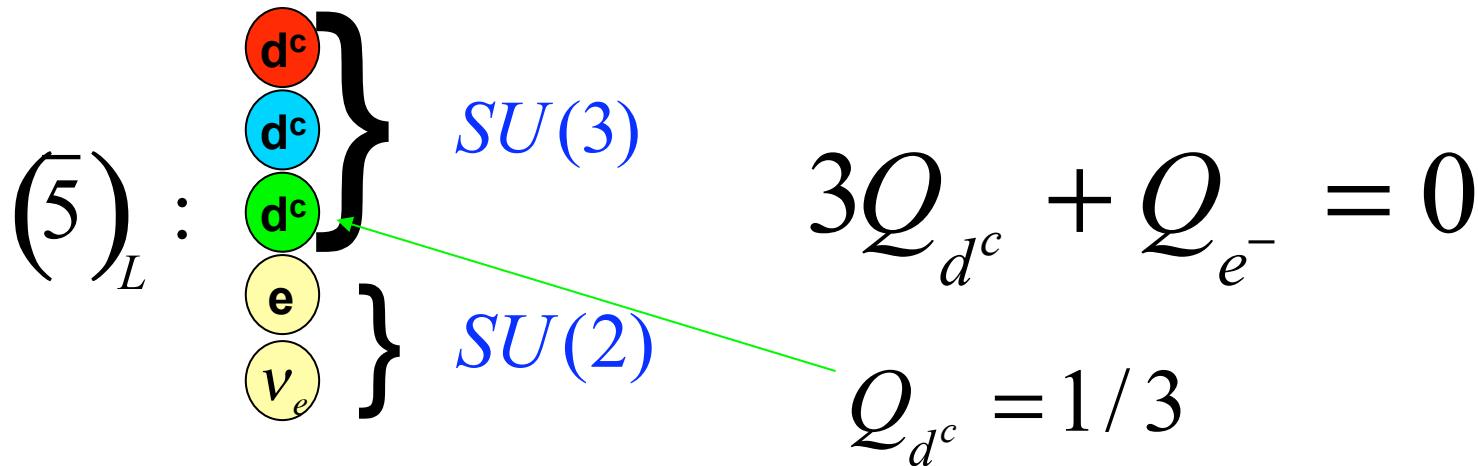
$$SU(5) \supset SU(3) \otimes SU(2) \otimes U(1)$$



$$3Q_{d^c} + Q_{e^-} = 0$$

Grand Unification

$$SO(10) \supset SU(5) \supset SU(3) \otimes SU(2) \otimes U(1)$$

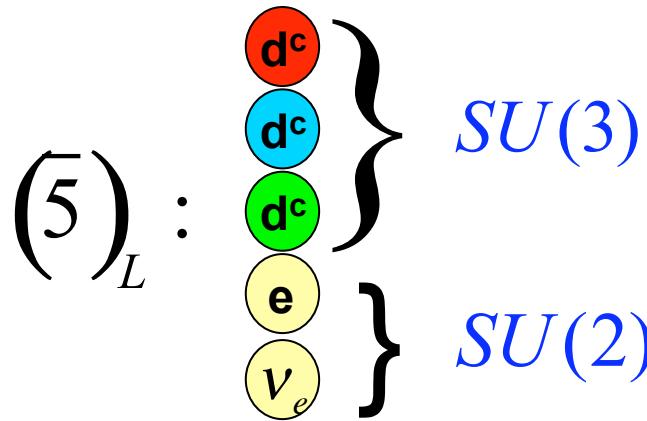


$$(16)_L = (10)_L + (\bar{5})_L + (1)_L$$

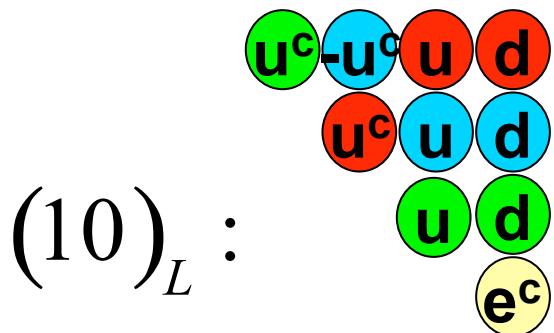
$$\nu_{e,L}^c \equiv \nu_{e,R}$$

Grand Unification

$$SO(10) \supset SU(5) \supset SU(3) \otimes SU(2) \otimes U(1)$$



$$3Q_{d^c} + Q_{e^-} = 0$$



$$(16)_L = (10)_L + (\bar{5})_L + (1)_L$$

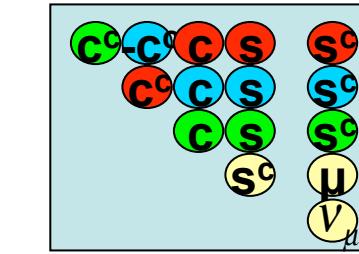
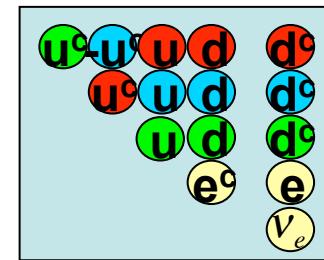
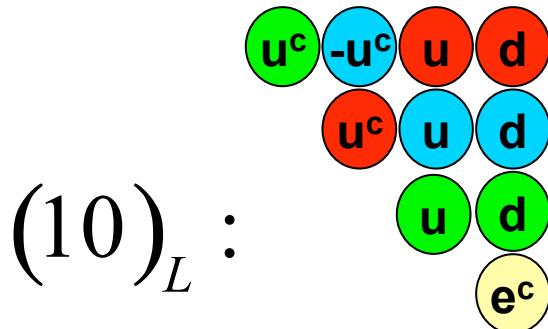
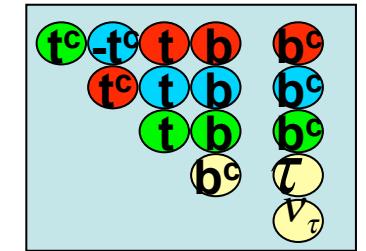
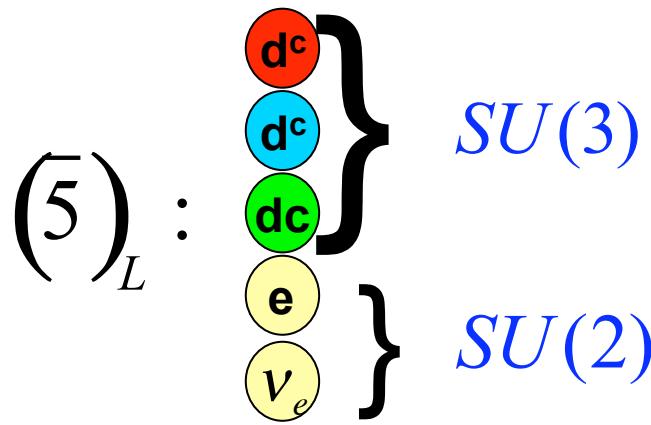
Neutrino masses via seesaw

The diagram illustrates the seesaw mechanism. It shows a horizontal line with a cross symbol in the center, representing the Majorana mass term. Two vertical lines branch off from the ends of the horizontal line, each ending in a vertex. From each vertex, a line goes up to a yellow circle labeled H and another line goes down to a yellow circle labeled v_L . A third line, labeled v_R , connects the two vertices.

$$M_{v_L} = \frac{\langle H \rangle^2}{M_{v_R}}$$

Grand Unification

$SO(10) \supset SU(5) \supset SU(3) \otimes SU(2) \otimes U(1)$



Generations

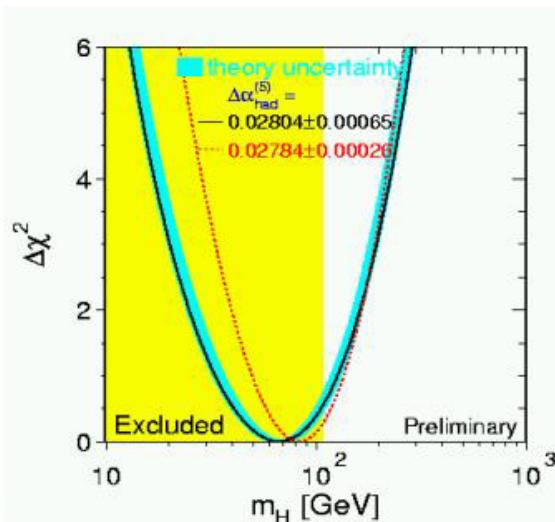
$$(16)_L = (10)_L + (\bar{5})_L + (1)_L$$

Precision tests

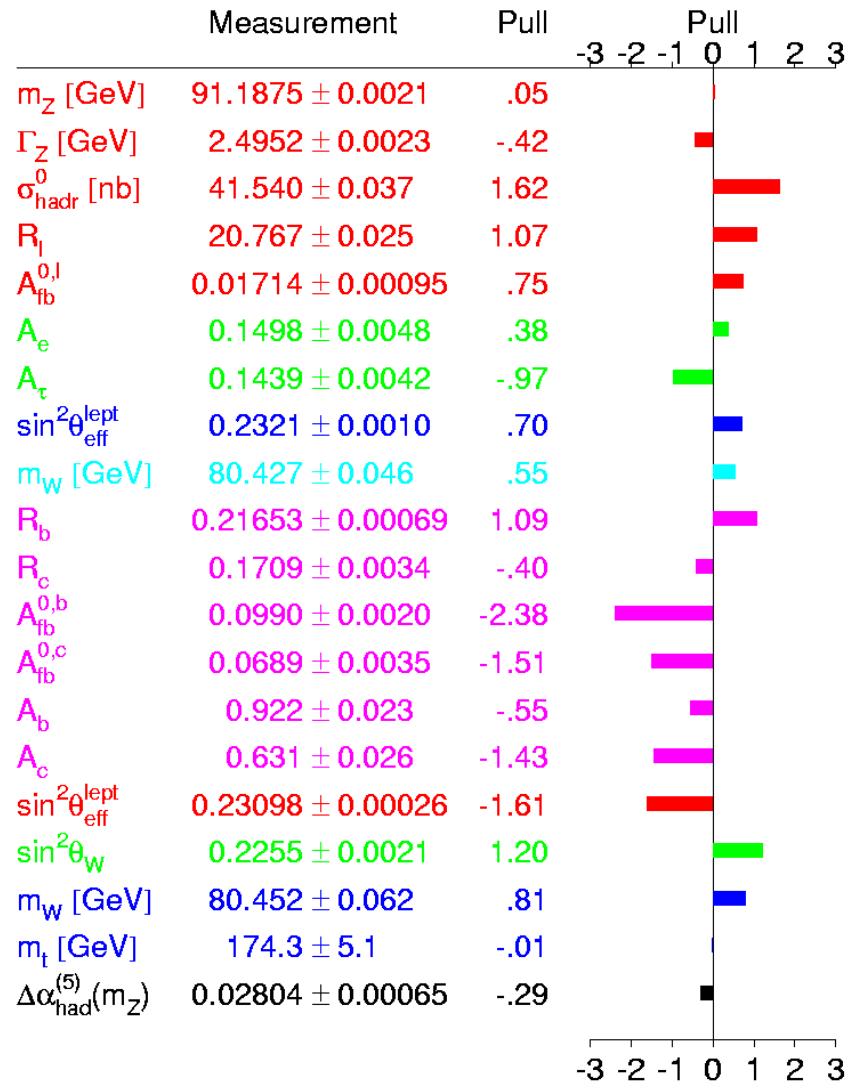
New families? :

$$N_\nu = 2.984 \pm 0.008$$

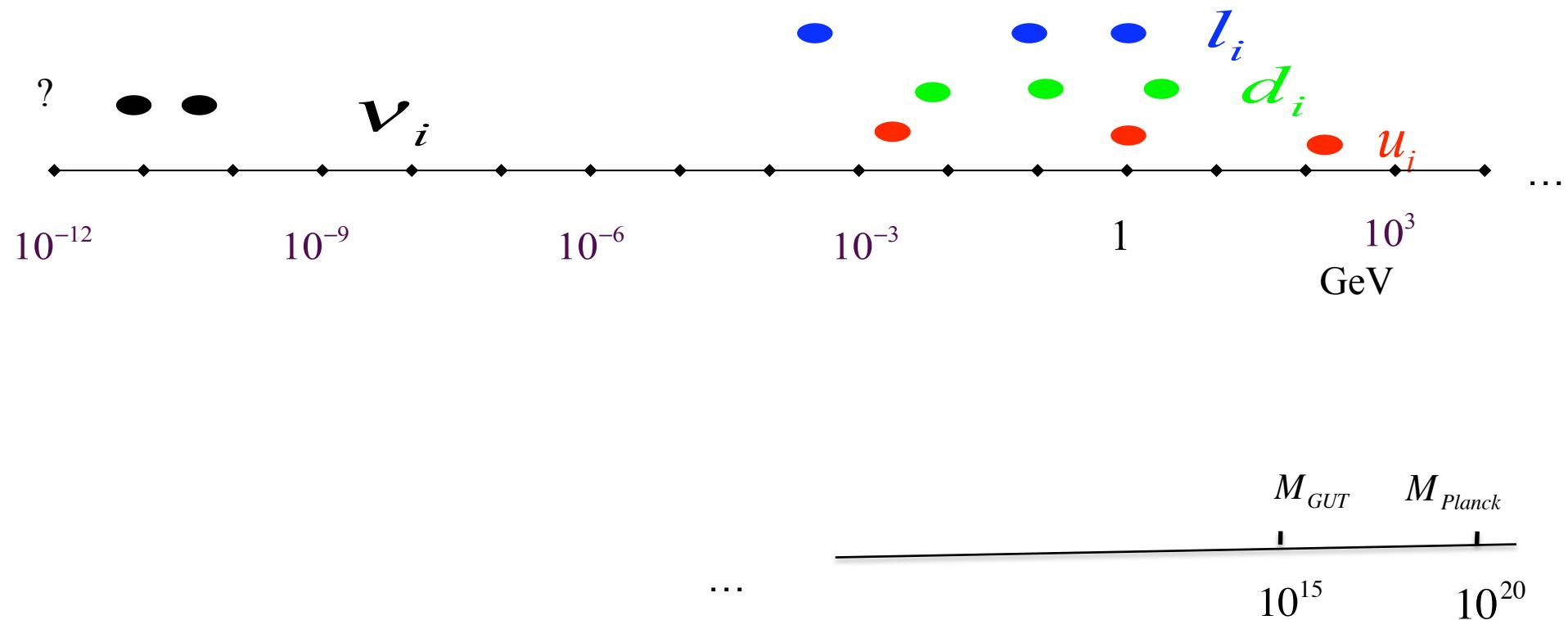
Higgs $\tilde{H} = (H^0 + \overline{H}^0) / \sqrt{2}$



$$\begin{array}{c} \dagger \\ \left(\begin{array}{c} H^0 \\ H^- \end{array} \right) \end{array} \left\{ W_\mu^\pm \right.$$



Mass scales



The Standard model as an effective field theory...

$$SU(3) \times SU(2) \times U(1) : G_\mu^{a=1..8}, W_\mu^{a=1..3}, B_\mu \quad A_\mu \rightarrow A_\mu + \partial_\mu \theta$$

$$\begin{pmatrix} u \\ d \end{pmatrix}_L, u_R, d_R, \begin{pmatrix} v \\ e \end{pmatrix}_L, e_R, v_R (?) \quad f_R \rightarrow e^{i\alpha_K} f_R$$

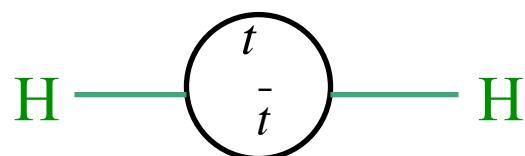
$$\begin{pmatrix} H^+ \\ H^0 \end{pmatrix}$$

$$L_{effective}^{SM} \supset \cancel{M_A} A_\mu A^\mu + \cancel{m_f} \overline{f_L} f_R + M_H^2 |H|^2$$

$$M_A, m_f \ll M_X, M_{Planck} ?$$

The hierarchy problem :

M_H not forbidden by SM symmetry:



$$M_H^2 \simeq \frac{h_t^2}{16\pi^2} \int_0^{\Lambda^2} dk^2 = \frac{h^2}{16\pi^2} \Lambda^2 \quad \Lambda \leq 1 TeV ??$$

The Standard model as an effective field theory...

A renormalisable, spontaneously broken, local gauge quantum field theory

- Renormalisable ✓
- Vectors gauge bosons ✓
- Fermions chiral ✓
- Massless gauge bosons - vectorlike couplings ✓
- Massive gauge bosons - chiral couplings ✓
- Spontaneously broken ✗ (hierarchy problem)

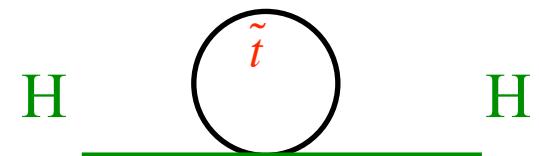
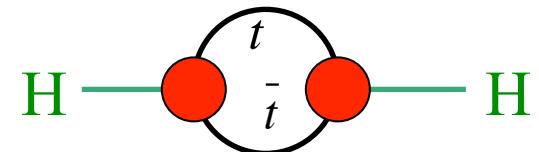
Solutions to the hierarchy problem

$\Lambda \leq 1TeV??$

- Composite: technicolour, walking technicolor, strongly coupled Standard Model,...

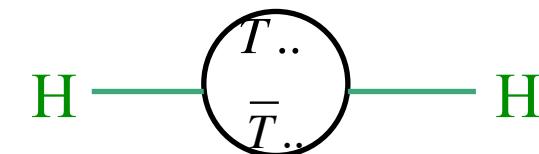
- Symmetry protection

SUSY



Goldstone: little Higgs

Double protection

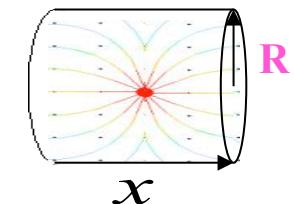


- $\Lambda_{fundamental} \simeq 1TeV!$

Xtra dimensions

$$V(r) = \frac{1}{M_*^{2+d} R^d} \frac{m_1 m_2}{r}, \quad D = 4 + d, \quad r \ll R$$

$$M_{Planck}^2 = M_*^2 (M_* R)^d$$



- (● Anthropic)

Split SUSY....