

# Relativistic theory of scattering

Scalar particle – satisfies KG equation

$$(\partial_\mu \partial^\mu + m^2)\phi = 0$$

$$\partial^\mu = \left(\frac{1}{c} \frac{\partial}{\partial t}, -\nabla\right), \quad \partial_\mu = \left(\frac{1}{c} \frac{\partial}{\partial t}, \nabla\right)$$

- Classical electrodynamics, motion of charge  $-e$  in EM potential  $A^\mu = (A^0, \mathbf{A})$  is obtained by the substitution :  $p^\mu \rightarrow p^\mu + eA^\mu$
- Quantum mechanics :  $i\partial^\mu \rightarrow i\partial^\mu + eA^\mu$

The Klein Gordon equation becomes:

$$(\partial_\mu \partial^\mu + m^2)\psi = -V\psi \quad \text{where} \quad V = -ie(\partial_\mu A^\mu + A^\mu \partial_\mu) - e^2 A^2$$

The smallness of the EM coupling,  $\alpha_{em} = \frac{e^2}{4\pi} \sim \frac{1}{137}$ , means that it is sensible to

Make a “perturbation” expansion of  $V$  in powers of  $\alpha_{em}$

Want to solve :

$$(\partial_\mu \partial^\mu + m^2)\psi = -V\psi$$

Solution :

$$\psi(x) = \phi(x) - \int d^4x' \Delta_F(x'-x)V(x')\psi(x')$$

where

$$(\partial_\mu \partial^\mu + m^2)\phi = 0$$

and

$$(\partial_\mu \partial^\mu + m^2)\Delta_F(x'-x) = \delta^4(x'-x)$$

Feynman propagator

Dirac Delta function

In bra- ket- notation

$$|\psi\rangle = |\phi\rangle - \frac{1}{\partial^2 + m^2 + i\epsilon} V |\psi\rangle$$

$$\psi(x) = \phi(x) - \int d^4x' \langle x | \frac{1}{\partial^2 + m^2 + i\epsilon} | x' \rangle \langle x' | V | \psi \rangle$$

# The propagator of the Klein Gordon particle

$$\begin{aligned}
 \Delta_F(x'-x) &= \langle x | \frac{1}{\partial^2 + m^2 + i\epsilon} | x' \rangle \\
 &= \int d^4 p' \int d^4 p'' \langle x | p' \rangle \underbrace{\langle p' | \frac{1}{\partial^2 + m^2 + i\epsilon} | p'' \rangle}_{\frac{\delta^4(p'-p'')}{-p'^2 + m^2 + i\epsilon}} \langle p'' | x' \rangle \\
 &= \int d^4 p' \left( \frac{1}{2\pi} \right)^2 e^{ip' \cdot x} \frac{1}{-p'^2 + m^2 + i\epsilon} \left( \frac{1}{2\pi} \right)^2 e^{-ip' \cdot x'}
 \end{aligned}$$

$$\Delta_F(x) = -\frac{1}{(2\pi)^4} \int d^4 p' e^{-ip' \cdot x} \frac{1}{p'^2 - m^2 - i\epsilon}$$

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Equivalent method : solve for propagator in momentum space by taking Fourier transform

$$\frac{1}{(2\pi)^2} \int e^{-ip \cdot (x'-x)} (\partial_\mu \partial^\mu + m^2)\Delta_F(x'-x)d^4(x'-x) = \frac{1}{(2\pi)^2} \int e^{-ip \cdot (x'-x)} \delta^4(x'-x)d^4(x'-x)$$

$$\Rightarrow (-p^2 + m^2)\Delta_F(p) = \frac{1}{(2\pi^2)}$$

$$\tilde{\Delta}_F(p) = \frac{1}{(2\pi)^2} \frac{1}{-p^2 + m^2 + i\epsilon}, \quad \Delta_F(x) = -\frac{1}{(2\pi)^4} \int d^4p e^{-ip \cdot x} \frac{1}{p^2 - m^2 - i\epsilon}$$

$$\Delta_F(x) = -\frac{1}{(2\pi)^4} \int d^4 p e^{-ip \cdot x} \frac{1}{p^2 - m^2 - i\epsilon}$$

$$p^2 + m^2 - i\epsilon \Rightarrow p_0^2 = \underline{p}^2 + m^2 - i\epsilon \Rightarrow p_0 = \pm \left( \underline{p}^2 + m^2 \right)^{1/2} \mp i\delta = \pm \omega_p \mp i\delta$$

$$\Delta_F(x' - x) = -\frac{1}{(2\pi)^4} \int d^3 p e^{-i\underline{p} \cdot (\underline{x}' - \underline{x})} \int dp_0 \frac{e^{-ip_0(t'-t)}}{\underbrace{(p_0 - (\omega_p - i\delta))(p_0 - (-\omega_p + i\delta))}_I}$$

- If  $t' - t > 0$ , choose contour such that  $p_0 = -ip_I$  ( $p_I + ve$ )  $\Rightarrow e^{-ip_0(t'-t)} = e^{-p_I(t'-t)}$

$$I = -\frac{\pi i}{\omega_p} e^{-i\omega_p(t'-t)} \theta(t'-t)$$

- If  $t' - t < 0$ , choose contour such that  $p_0 = +ip_I$  ( $p_I + ve$ )  $\Rightarrow e^{-ip_0(t'-t)} = e^{-p_I(t'-t)}$

$$I = -\frac{\pi i}{\omega_p} e^{+i\omega_p(t'-t)} \theta(t-t')$$

# Feynman propagator

$$\Delta_F(x-x') = \frac{i}{(2\pi)^3 2\omega_p} \int d^3 p \left[ e^{-i\omega_p(t'-t) + i\underline{p}\cdot(\underline{x}'-\underline{x})} \theta(t'-t) + e^{+i\omega_p(t'-t) + i\underline{p}\cdot(\underline{x}'-\underline{x})} \theta(t-t') \right]$$

An alternative form :

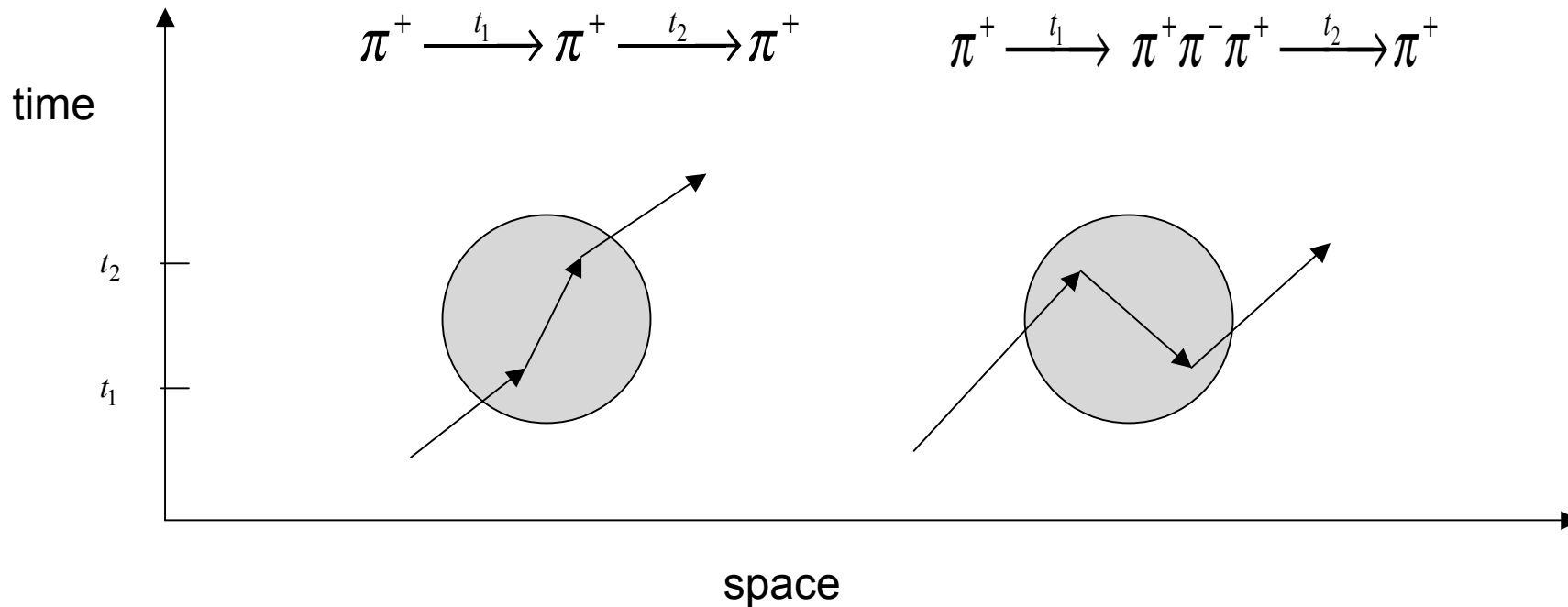
$$f_p^\pm = e^{\mp i p \cdot x} \frac{1}{\sqrt{2 p^0 V}} = e^{\mp i p \cdot x} \frac{1}{\sqrt{2 p^0} (2\pi)^3}$$

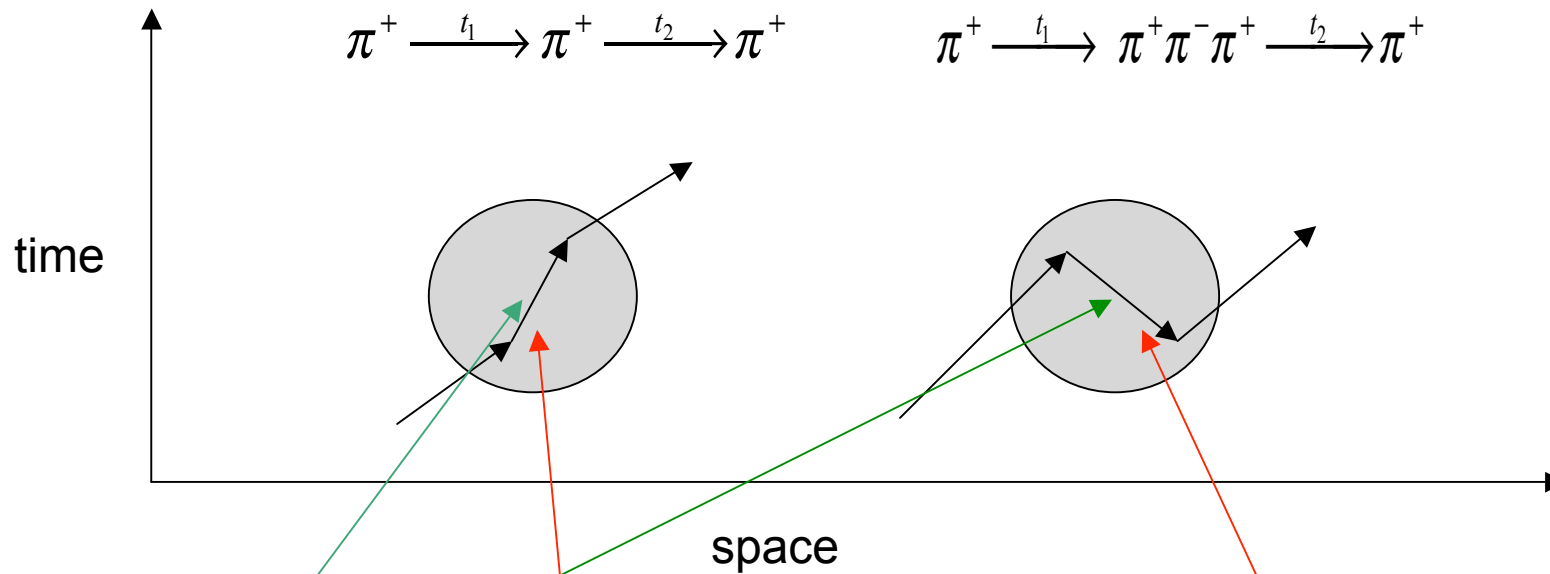
$$\Delta_F(x-x') = i \int d^3 p f_p^+(x') f_p^{+*}(x) \theta(t'-t) + i \int d^3 p f_p^-(x') f_p^{-*}(x) \theta(t-t')$$

# Feynman – Stuckelberg interpretation

$$\begin{array}{ccc}
 \pi^+(E > 0) & \uparrow & \equiv & \pi^-(E < 0) & \downarrow \\
 & e^{-iEt} & & & e^{-i(-E)(-t)}
 \end{array}$$

Two different time orderings giving same observable event :





$$\Delta_F(x'-x) = \frac{1}{(2\pi)^4} \int d^4 p e^{-ip \cdot (x'-x)} \frac{1}{p^2 - m^2 - i\epsilon} = i \int \frac{d^3 p}{(2\pi)^3 2p^0} e^{-ip^0 |t'-t| - ip \cdot (x'-x)}$$

$$\Delta_F(x-x') = i \int d^3 p f_p^+(x') f_p^{+*}(x) \theta(t_2 - t_1) + i \int d^3 p f_p^-(x') f_p^{-*}(x) \theta(t_1 - t_2)$$

where  $f_p^\pm = e^{\mp i p \cdot x} \frac{1}{\sqrt{2p^0 V}}$  are positive and negative energy solutions to free KG equation