

PHYS 543 Fall 2017

Homework 1

Solutions

1.1 a) $\mathcal{N}^2 (\langle \alpha | a \rangle (a^\dagger | \alpha \rangle) = 1$

$$\mathcal{N}^2 \langle \alpha | (a^\dagger a + 1) | \alpha \rangle = 1$$

$$\mathcal{N}^2 (|\alpha|^2 + 1) = 1$$

$$\mathcal{N} = \frac{1}{\sqrt{|\alpha|^2 + 1}}$$

b) $|\alpha\rangle = e^{-|\alpha|^2/2} \sum_n \frac{\alpha^n}{\sqrt{n!}} |n\rangle$

$$\mathcal{N} a^\dagger |\alpha\rangle = \mathcal{N} e^{-|\alpha|^2/2} \sum_n \frac{\alpha^{n+1}}{\sqrt{n!}} |n+1\rangle$$

c) $\langle x \rangle = \mathcal{N}^2 \langle \alpha | \hat{x} | \alpha \rangle = \frac{1}{\sqrt{2}} \mathcal{N}^2 \langle \alpha | a (a + a^\dagger) a^\dagger | \alpha \rangle$

$$\begin{aligned} a(a+a^\dagger)a^\dagger &= a a a^\dagger + a a^\dagger a^\dagger \\ &= a(a^\dagger a + 1) + (a^\dagger a + 1)a^\dagger \\ &= a a^\dagger a + a + a^\dagger a a^\dagger + a^\dagger \\ &= a^\dagger a a + 2a + a^\dagger a^\dagger a + 2a^\dagger \end{aligned}$$

$$\begin{aligned} \Rightarrow \langle x \rangle &= \frac{1}{\sqrt{2}} \mathcal{N}^2 (|\alpha|^2 \alpha + 2\alpha + |\alpha| |\alpha|^* + 2\alpha^*) \\ &= \frac{1}{\sqrt{2}} \frac{(|\alpha|^2 + 2)(\alpha + \alpha^*)}{|\alpha|^2 + 1} = \frac{|\alpha|^2 + 2}{|\alpha|^2 + 1} \text{Re } \alpha \sqrt{2} \end{aligned}$$

d) $\langle x | \alpha \rangle = \Psi_0(x - \alpha\sqrt{2}) = \frac{1}{\pi^{1/4}} e^{-(x - \alpha\sqrt{2})^2/2}$

$$\begin{aligned} \mathcal{N} \langle x | a^\dagger | \alpha \rangle &= \mathcal{N} \langle x | \frac{x - iP}{\sqrt{2}} | \alpha \rangle \\ &= \frac{\mathcal{N}}{\sqrt{2}} (x - \frac{d}{dx}) e^{-(x - \alpha\sqrt{2})^2/2} \end{aligned}$$

$$= \frac{1}{\sqrt{2} \sqrt{2} (|\alpha|^2 + 1)} (2x - \alpha\sqrt{2}) e^{-(x - \alpha\sqrt{2})^2/2}$$

e) $d \rightarrow 0 \quad a^\dagger |d\rangle \rightarrow |1\rangle$

$d \rightarrow \infty \quad a^\dagger |d\rangle \rightarrow |d\rangle$

$$f) |\psi\rangle = D(c|0\rangle + d|1\rangle)$$

$$(\text{denote } \hat{D}_x(\alpha\sqrt{2}) \equiv \hat{D})$$

$$= c\hat{D}|0\rangle + d\hat{D}\hat{a}^\dagger|0\rangle$$

$$(\text{use } \hat{a}^\dagger\hat{D} = \hat{D}\hat{a}^\dagger + \alpha\hat{D})$$

$$= c\hat{D}|0\rangle + d\hat{a}^\dagger\hat{D}|0\rangle - \alpha d\hat{D}|0\rangle$$

$$= (c - \alpha d)|0\rangle + d\hat{a}^\dagger|0\rangle$$

$$= d\hat{a}^\dagger|0\rangle \quad \text{if } c - \alpha d = 0$$

\Downarrow

$$c = \frac{-\alpha}{\sqrt{1+\alpha^2}}, \quad d = \frac{1}{\sqrt{1+\alpha^2}} \quad (\text{normalized})$$

$$\text{Then } |\psi\rangle = \frac{\hat{a}^\dagger|0\rangle}{\sqrt{1+\alpha^2}} \quad \text{consistent with (a)}$$

$$\langle x|\psi\rangle = c\psi_0(x - \alpha\sqrt{2}) + d\psi_1(x - \alpha\sqrt{2})$$

$$= \frac{1}{\pi^{1/4}\sqrt{1+\alpha^2}} e^{-(x-\alpha\sqrt{2})^2} (+\alpha + \sqrt{2}(x-\alpha\sqrt{2}))$$

$$= \frac{1}{\pi^{1/4}\sqrt{1+\alpha^2}} e^{-(x-\alpha\sqrt{2})^2} (x\sqrt{2} - \alpha)$$

consistent with (d).

1.2

$$\begin{aligned}
 \text{a) } [X_\theta, X_{\frac{\pi}{2}+\theta}] &= [X_C + P_S, -X_S + P_C] \\
 &= c^2[X, P] - s^2[P, X] = i \\
 \langle \Delta X_\theta^2 \rangle &\langle \Delta X_{\frac{\pi}{2}+\theta}^2 \rangle \geq \frac{1}{4}
 \end{aligned}$$

$$\begin{aligned}
 \text{b) For } \theta=0: H &= -\frac{\hbar\gamma}{2} [XP + PX] \\
 \Rightarrow \text{for arbitrary } \theta:
 \end{aligned}$$

$$H = -\frac{\hbar\gamma}{2} [X_\theta X_{\frac{\pi}{2}+\theta} + X_{\frac{\pi}{2}+\theta} X_\theta] \quad (1)$$

$$\begin{aligned}
 \text{Then } \dot{X}_\theta &= \frac{i}{\hbar} [H, X_\theta] = \\
 &= \frac{-i\gamma}{2} (X_\theta [X_{\frac{\pi}{2}+\theta}, X_\theta] + [X_{\frac{\pi}{2}+\theta}, X_\theta] X_\theta) \\
 &= -\gamma X_\theta
 \end{aligned}$$

$$\dot{X}_{\frac{\pi}{2}+\theta} = \frac{i}{\hbar} [H, X_{\frac{\pi}{2}+\theta}] = \gamma X_{\frac{\pi}{2}+\theta}$$

$$\begin{aligned}
 \text{c) } X_\theta &= \frac{a+a^\dagger}{\sqrt{2}} \frac{e^{i\theta}+e^{-i\theta}}{2} + \frac{a-a^\dagger}{\sqrt{2}i} \frac{e^{i\theta}-e^{-i\theta}}{2i} \\
 &= \frac{1}{\sqrt{2}} (a e^{-i\theta} + a^\dagger e^{i\theta})
 \end{aligned}$$

$$\text{From (1): } \hat{H} = -\frac{\hbar\gamma}{2} (X_\theta X_{\frac{\pi}{2}+\theta} - i)$$

$$= \frac{-\hbar\gamma}{2} [(a e^{-i\theta} + a^\dagger e^{i\theta})(-i a e^{-i\theta} + i a^\dagger e^{i\theta}) - i]$$

$$= \frac{-\hbar\gamma}{2} (-i a^2 e^{-2i\theta} + i a^{\dagger 2} e^{2i\theta} - \cancel{i a^\dagger a} + \cancel{i a a^\dagger} - i)$$

$$e^{-\frac{i}{\hbar} \hat{H} t} |0\rangle = \left(\hat{1} - \frac{i}{\hbar} \hat{H} t \right) |0\rangle$$

$$= |0\rangle - \frac{\gamma t}{2} a^{\dagger 2} e^{2i\theta} |0\rangle$$

$$= |0\rangle - \frac{\gamma t}{\sqrt{2}} e^{2i\theta} |2\rangle$$

$$d) \langle \Psi | X_{\theta} | \Psi \rangle = \langle \Psi | X_{\frac{\pi}{2}+\theta} | \Psi \rangle = 0$$

because $X_{\theta} | \Psi \rangle$ and $X_{\frac{\pi}{2}+\theta} | \Psi \rangle$ only contain odd photon-number terms.

$$\langle \Delta X_{\theta}^2 \rangle = \langle \Psi | X_{\theta}^2 | \Psi \rangle$$

$$= \left(\langle 0 | -\frac{\gamma t}{\sqrt{2}} e^{-2i\theta} \langle 2 | \right) \left(\frac{a^2 e^{-2i\theta} + a a^\dagger + a^\dagger a + a^{\dagger 2} e^{2i\theta}}{2} \right) \left(| 0 \rangle - \frac{\gamma t}{\sqrt{2}} e^{2i\theta} | 2 \rangle \right)$$

$$\stackrel{\gamma t \ll 1}{=} \langle 0 | \frac{a a^\dagger}{2} | 0 \rangle + \frac{1}{2} \langle 0 | a^2 e^{-2i\theta} \left(\frac{\gamma t}{\sqrt{2}} e^{2i\theta} \right) | 2 \rangle$$

$$+ \frac{1}{2} \left(-\frac{\gamma t}{\sqrt{2}} e^{-2i\theta} \right) \langle 2 | a^{\dagger 2} e^{2i\theta} | 0 \rangle$$

$$= \frac{1}{2} (1 - 2\gamma t) \approx \frac{1}{2} e^{-2\gamma t}$$

Similar calculation for

$$\langle \Delta X_{\frac{\pi}{2}+\theta}^2 \rangle = \frac{1}{2} (1 + 2\gamma t) \approx \frac{1}{2} e^{2\gamma t}$$

$$e) H = \hbar \gamma (X_{\theta,A} X_{\frac{\pi}{2}+\theta,B} + X_{\frac{\pi}{2}+\theta,A} X_{\theta,B})$$

$$\dot{X}_{\theta,\pm} = \frac{i}{\hbar} [H, \frac{1}{\sqrt{2}} (X_{\theta,A} \pm X_{\theta,B})]$$

$$= \frac{i}{\hbar} \frac{\hbar \gamma}{\sqrt{2}} (X_{\theta,A} [X_{\frac{\pi}{2}+\theta,B}, \pm X_{\theta,B}] + [X_{\frac{\pi}{2}+\theta,A}, X_{\theta,A}] X_{\theta,B})$$

$$= \frac{i\gamma}{\sqrt{2}} (\mp i X_{\theta,A} - i X_{\theta,B}) = \pm \gamma X_{\theta,\pm}$$

$$\dot{X}_{\frac{\pi}{2}+\theta,\pm} = \frac{i}{\hbar} [H, \frac{1}{\sqrt{2}} (X_{\frac{\pi}{2}+\theta,A} \pm X_{\frac{\pi}{2}+\theta,B})]$$

$$= \frac{i\gamma}{\sqrt{2}} (i X_{\frac{\pi}{2}+\theta,B} \pm i X_{\frac{\pi}{2}+\theta,A}) = \mp \gamma X_{\frac{\pi}{2}+\theta,\pm}$$

$$d) H = \hbar \gamma \left(\frac{a_A e^{-i\theta} + a_A^\dagger e^{i\theta}}{\sqrt{2}} \cdot \frac{-i a_B e^{-i\theta} + i a_B^\dagger e^{i\theta}}{\sqrt{2}} + \frac{-i a_A e^{-i\theta} + i a_A^\dagger e^{i\theta}}{\sqrt{2}} \cdot \frac{a_B e^{-i\theta} + a_B^\dagger e^{i\theta}}{\sqrt{2}} \right)$$

$$= i \hbar \gamma \left(-a_A a_B e^{-2i\theta} + a_A^\dagger a_B^\dagger e^{2i\theta} \right)$$

$$e^{-\frac{i}{\hbar} H t} |00\rangle = |00\rangle + \gamma t a_A^\dagger a_B^\dagger e^{2i\theta} |00\rangle$$

$$= |00\rangle + \gamma t e^{2i\theta} |11\rangle = |\Psi\rangle$$

g) Following Ex. 3.113:

$$\langle \Psi | X_{\theta, \pm} | \Psi \rangle = 0$$

$$\langle \Delta X_{\theta, \pm}^2 \rangle = \langle \Psi | X_{\theta, \pm}^2 | \Psi \rangle$$

$$= \left(\langle 00 | + \gamma t e^{2i\theta} \langle 11 | \right) \frac{a_A a_A^\dagger + a_A^\dagger a_A + a_B a_B^\dagger + a_B^\dagger a_B \pm 2 a_A a_B e^{-2i\theta} \pm 2 a_A^\dagger a_B^\dagger e^{2i\theta}}{4}$$

$$\times \left(|00\rangle + \gamma t e^{2i\theta} |11\rangle \right)$$

$$= \langle 00 | \frac{a_A a_A^\dagger + a_B a_B^\dagger}{4} |00\rangle \pm \gamma t e^{-2i\theta} \langle 11 | \frac{a_A^\dagger a_B^\dagger e^{2i\theta}}{2} |00\rangle$$

$$\pm \gamma t e^{2i\theta} \langle 00 | \frac{a_A a_B e^{-2i\theta}}{2} |11\rangle$$

$$= \frac{1}{2} \pm \gamma t \approx \frac{1}{2} e^{\pm 2\gamma t}$$

Similarly,

$$\langle \Delta X_{\frac{\pi}{2} + \theta, \pm} \rangle = \frac{1}{2} \pm \gamma t \approx \frac{1}{2} e^{\mp 2\gamma t}$$

$$\begin{aligned}
 \boxed{1.3} \quad [L_j, [L_k, L_e]] &= [L_j, i\hbar \sum_{m,n} \epsilon_{kem} L_m] \\
 &= (i\hbar)^2 \sum_{k,e,m} \epsilon_{jem} L_m \\
 &= -\hbar^2 \sum_{m,k,e} \epsilon_{mke} \epsilon_{muj} L_m \\
 &= \hbar^2 L_m (-\delta_{km} \delta_{ej} + \delta_{kj} \delta_{em}) \quad (1)
 \end{aligned}$$

$$\begin{aligned}
 [L_x, [L_y, L_z]] &= [L_x, i\hbar L_x] = 0 \\
 j=1, \quad k=2, \quad e=3 &\Rightarrow \text{all } \delta\text{'s vanish in (1)}
 \end{aligned}$$

$$[L_x, [L_x, L_z]] = [L_x, -i\hbar L_y] = \hbar^2 L_z$$

$j=k=1, \quad e=3 \Rightarrow$ second term in (1) nonzero for $m=3$

$$[L_x, [L_z, L_x]] = [L_x, i\hbar L_y] = -\hbar^2 L_z$$

$j=e=1, \quad k=3 \Rightarrow$ first term in (1) nonzero for $m=3$.

1.4

$$x = r \sin \theta \cos \varphi$$

$$y = r \sin \theta \sin \varphi$$

$$z = r \cos \theta$$

$$\frac{\partial}{\partial \theta} = \frac{\partial}{\partial x} \frac{\partial x}{\partial \theta} + \frac{\partial}{\partial y} \frac{\partial y}{\partial \theta} + \frac{\partial}{\partial z} \frac{\partial z}{\partial \theta}$$

$$= r \cos \theta \cos \varphi \frac{\partial}{\partial x} + r \cos \theta \sin \varphi \frac{\partial}{\partial y} - r \sin \theta \frac{\partial}{\partial z}$$

$$\frac{\partial}{\partial \varphi} = \frac{\partial}{\partial x} \frac{\partial x}{\partial \varphi} + \frac{\partial}{\partial y} \frac{\partial y}{\partial \varphi} + \frac{\partial}{\partial z} \frac{\partial z}{\partial \varphi}$$

$$= -r \sin \theta \sin \varphi \frac{\partial}{\partial x} + r \sin \theta \cos \varphi \frac{\partial}{\partial y}$$

$$L_x = i\hbar \left(\sin \varphi \frac{\partial}{\partial \theta} + \cot \theta \cos \varphi \frac{\partial}{\partial \varphi} \right)$$

$$= i\hbar \left(r \cos \theta \sin \varphi \cos \varphi \frac{\partial}{\partial x} + r \cos \theta \sin^2 \varphi \frac{\partial}{\partial y} - r \sin \theta \sin \varphi \frac{\partial}{\partial z} \right.$$

$$\left. - r \frac{\cos \theta}{\sin \theta} \cos \varphi \sin \theta \sin \varphi \frac{\partial}{\partial x} + r \frac{\cos \theta}{\sin \theta} \cos^2 \varphi \sin \theta \frac{\partial}{\partial y} \right)$$

$$= i\hbar \left(r \cos \theta \frac{\partial}{\partial y} - r \sin \theta \sin \varphi \frac{\partial}{\partial z} \right)$$

$$= i\hbar \left(z \frac{\partial}{\partial y} - y \frac{\partial}{\partial z} \right) = y \hat{p}_z - z \hat{p}_y$$

$$L_y = i\hbar \left(-\cos \varphi \frac{\partial}{\partial \theta} + \cot \theta \sin \varphi \frac{\partial}{\partial \varphi} \right)$$

$$= i\hbar \left(-r \cos \theta \cos^2 \varphi \frac{\partial}{\partial x} - r \cos \theta \sin \varphi \cos \varphi \frac{\partial}{\partial y} + r \sin \theta \cos \varphi \frac{\partial}{\partial z} \right.$$

$$\left. - r \frac{\cos \theta}{\sin \theta} \sin^2 \varphi \sin \theta \frac{\partial}{\partial x} + r \frac{\cos \theta}{\sin \theta} \sin \theta \sin \varphi \cos \varphi \frac{\partial}{\partial y} \right)$$

$$= i\hbar \left(-r \cos \theta \frac{\partial}{\partial x} + r \sin \theta \cos \varphi \frac{\partial}{\partial z} \right)$$

$$= i\hbar \left(-z \frac{\partial}{\partial x} + x \frac{\partial}{\partial z} \right) = z \hat{p}_x - x \hat{p}_z$$

$$L_z = -i\hbar \frac{\partial}{\partial \varphi} = -i\hbar \left(-y \frac{\partial}{\partial x} + x \frac{\partial}{\partial y} \right) = x \hat{p}_y - y \hat{p}_x$$