

Solutions

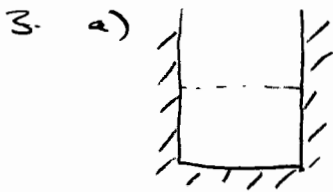
1. a) No: $\hat{R} \rightarrow \begin{pmatrix} 0 & 0 & 1 \\ 1 & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix}$ while $\hat{R}^\dagger \rightarrow \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{pmatrix}$ so $\hat{R} \neq \hat{R}^\dagger$

b) $\hat{H} = \hbar\omega \begin{pmatrix} 0 & 1 & 1 \\ 1 & 0 & 1 \\ 1 & 1 & 0 \end{pmatrix}$ Characteristic equation: $\det \begin{pmatrix} -\lambda & 1 & 1 \\ 1 & -\lambda & 1 \\ 1 & 1 & -\lambda \end{pmatrix} = 0$
 $\lambda^3 - 3\lambda - 2 = 0$
 $(\lambda + 1)^2 (\lambda - 2) = 0$

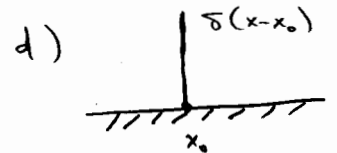
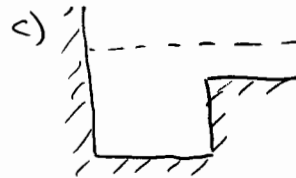
Eigenvalue $\lambda = -1$ ($\times \hbar\omega$): eigenstates $\frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \\ -1 \end{pmatrix} \rightarrow \frac{1}{\sqrt{2}} (|B\rangle - |C\rangle)$

$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 0 \\ -1 \end{pmatrix} \rightarrow \frac{1}{\sqrt{2}} (|A\rangle - |C\rangle)$

Eigenvalue $\lambda = 2$ ($\times \hbar\omega$): eigenstate $\frac{1}{\sqrt{3}} \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} \rightarrow \frac{1}{\sqrt{3}} (|A\rangle + |B\rangle + |C\rangle)$



b) none (discontinuity)



In (d), the delta-function is positive because the second derivative of $V(x)$ at x_0 is $+\infty$. The wave function is stationary if there are potential barriers somewhere to left and right.

2. Ex. 3.43: $\psi_{t=0}(x) = \sqrt{k_0} \begin{cases} e^{-k_0 x} & x > 0 \\ e^{k_0 x} & x < 0 \end{cases} \quad k_0 = \frac{mV_0}{\hbar^2}$

$$\psi_{t>0}(x) = \sqrt{2k_0} \begin{cases} e^{-2k_0 x} \\ e^{2k_0 x} \end{cases}$$

(bound stationary solutions of the Schrödinger equation)

The probability that a bound state at $t=0$ will remain bound is

$$P = \left| \int \psi_{t=0}(x) \psi_{t>0}(x) dx \right|^2 = \left| 2\sqrt{2} k_0 \int_0^\infty e^{-3k_0 x} dx \right|^2 = \left| \frac{2\sqrt{2}}{3} \right|^2 = \frac{8}{9}$$

If the state does not remain bound, it will be localized at infinity at $t = \infty$. The probability to detect a finite position value is $P = 8/9$

$$4. \hat{H} = J \sum_A \sum_B (\hat{S}_{xA} \hat{S}_{xB} + \hat{S}_{yA} \hat{S}_{yB} + \hat{S}_{zA} \hat{S}_{zB})$$

$$S_x \rightarrow \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$S_{xA} S_{xB} \rightarrow \frac{\hbar^2}{4} \begin{pmatrix} (\frac{1}{2} \ \frac{1}{2}) & (\frac{1}{2} \ -\frac{1}{2}) & (-\frac{1}{2} \ \frac{1}{2}) & (-\frac{1}{2} \ -\frac{1}{2}) \\ (\frac{1}{2} \ \frac{1}{2}) & & & 1 \\ (\frac{1}{2} \ -\frac{1}{2}) & & 1 & \\ (-\frac{1}{2} \ \frac{1}{2}) & & & \\ (-\frac{1}{2} \ -\frac{1}{2}) & 1 & & \end{pmatrix}$$

$$S_y \rightarrow \frac{\hbar}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$S_{yA} S_{yB} \rightarrow \frac{\hbar^2}{4} \begin{pmatrix} & & & -1 \\ & & & \\ & & & \\ -1 & & & \end{pmatrix}$$

$$S_z \rightarrow \frac{\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$S_{zA} S_{zB} \rightarrow \frac{\hbar^2}{4} \begin{pmatrix} & & & \\ & & & \\ & & -1 & \\ & & & -1 \end{pmatrix}$$

$$\hat{H} \rightarrow \frac{J}{4} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 2 & 0 \\ 0 & 2 & -1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

Eigenvalues ($\times \frac{J\hbar^2}{4}$): (1, 1, 1, -3)

Eigenstates: $\begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix}$, $\begin{pmatrix} 0 \\ 0 \\ 0 \\ 1 \end{pmatrix}$, $\frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \\ -1 \\ 0 \end{pmatrix}$, $\frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \\ 1 \\ 0 \end{pmatrix}$

$$e^{-i\hat{H}t/\hbar} = e^{-\frac{iJ\hbar t}{4}} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} + e^{-\frac{iJ\hbar t}{4}} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} + e^{-\frac{iJ\hbar t}{4}} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 1 & 1 & 0 \\ 0 & 1 & 1 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} + e^{+\frac{3iJ\hbar t}{4}} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 1 & -1 & 0 \\ 0 & -1 & 1 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

For $J\hbar t = \pi$:

$$e^{-i\hat{H}t/\hbar} = e^{-i\pi/4} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \equiv \hat{U}$$

$$\text{If } |\psi\rangle = a|\uparrow\rangle + b|\downarrow\rangle$$

$$|\psi\rangle = c|\uparrow\rangle + d|\downarrow\rangle$$

(where $|\uparrow\rangle = |+\frac{1}{2}\rangle$, $|\downarrow\rangle = |-\frac{1}{2}\rangle$)

$$|\psi\rangle_A \otimes |\psi\rangle_B = ac|\uparrow\uparrow\rangle + ad|\uparrow\downarrow\rangle + bc|\downarrow\uparrow\rangle + bd|\downarrow\downarrow\rangle$$

$$\hat{U}(|\psi\rangle_A \otimes |\psi\rangle_B) = ac|\uparrow\uparrow\rangle + ad|\downarrow\uparrow\rangle + bc|\uparrow\downarrow\rangle + bd|\downarrow\downarrow\rangle = \\ = |\psi\rangle_A \otimes |\psi\rangle_B$$