

Classical Mechanics

LECTURE 12:

MORE ON ROCKETS;

NON-INERTIAL FRAMES

Prof. N. Harnew
University of Oxford
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OUTLINE : 12. MORE ON ROCKETS; NON-INERTIAL FRAMES

12.1 The rocket : vertical launch

12.2 The 1-stage vs. 2-stage rocket

12.3 Non-inertial reference frames

12.3.1 Commonplace examples

12.3.2 Example- accelerating lift

12.1 The rocket : vertical launch

- Rocket equation:

$$m \frac{dv}{dt} + w \frac{dm}{dt} = F$$

- Rocket rises against gravity

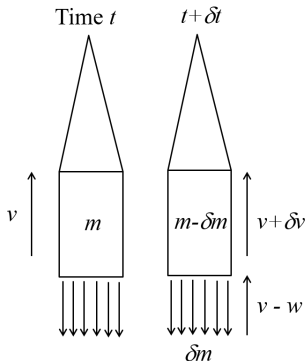
$$F = -mg$$

- Mass is ejected at *constant* velocity w relative to the rocket

- Rocket ejects mass uniformly:

$$m = m_0 - \alpha t$$

$$\rightarrow \frac{dm}{dt} = -\alpha$$



- Now consider upward motion:

$$m dv = (-mg + w\alpha) dt \rightarrow \int_{v_i}^{v_f} dv = \int_{t_i}^{t_f} \left(-g + \frac{w\alpha}{m_0 - \alpha t} \right) dt$$

$$\begin{aligned} v_f - v_i &= \left[-g(t_f - t_i) - w \log_e \frac{(m_0 - \alpha t_f)}{(m_0 - \alpha t_i)} \right] \\ &= \left[-g(t_f - t_i) - w \log_e (m_f / m_i) \right] \end{aligned}$$

Rocket vertical launch, continued

The rocket starts from rest at $t = 0$; half the mass is fuel. What is the velocity and height reached by the rocket at burn-out at time $t = T$?

$$\blacktriangleright v = \left[-gt - w \log_e \frac{(m_0 - \alpha t)}{(m_0)} \right] = \left[-gt - w \log_e \left(1 - \frac{\alpha}{m_0} t \right) \right] = \frac{dx}{dt}$$

$$\blacktriangleright \text{What is the condition for the rocket to rise ? } \rightarrow \frac{dv}{dt} > 0$$

$$\text{At } t = 0, m = m_0, \frac{dm}{dt} = -\alpha : \alpha w - m_0 g > 0 \rightarrow w > \frac{m_0 g}{\alpha}$$

$$\blacktriangleright m = m_0 - \alpha t ; \text{ at burnout } t = T, m = \frac{m_0}{2} \rightarrow \alpha = \frac{m_0}{2T}$$

\blacktriangleright Maximum velocity is at the burn-out of the fuel:

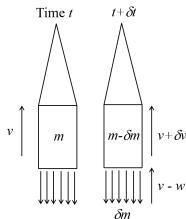
$$\text{At } t = T : v_{max} = -gT + w \log_e 2$$

$$\text{Height : } \int_0^x dx = \int_0^T \left[-gt - w \log_e \left(1 - \frac{\alpha}{m_0} t \right) \right] dt$$

$$\blacktriangleright \text{Standard integral : } \int \log_e z dz = z \log_e z - z$$

$$\blacktriangleright x = -\frac{gT^2}{2} + \frac{w m_0}{\alpha} \left[\left(1 - \frac{\alpha}{m_0} t \right) \left(\log_e \left(1 - \frac{\alpha}{m_0} t \right) \right) - \left(1 - \frac{\alpha}{m_0} t \right) \right]_0^T$$

$$\blacktriangleright \text{After simplification : } x = -\frac{gT^2}{2} + wT(1 - \log_e 2)$$



12.2 The 1-stage vs. 2-stage rocket

A two stage rocket is launched vertically from earth, total mass $M_0 = 10000$ kg and carries an additional payload of $m = 100$ kg. The fuel is 75% of the mass in both stages; burn rate $\alpha = 500$ kg s^{-1} , thrust velocity $w = 2.5$ km s^{-1} . The mass of the 2nd stage is 900 kg.

i) Calculate the final speed for the equivalent single stage rocket

ii) Find the final speed of the 2-stage rocket

(i) Single stage rocket :

- ▶ Time to burn-out :

$$\left| \frac{\Delta m}{\Delta t} \right| = \alpha \rightarrow T = 0.75(M_1 + M_2)/500 \text{ [kg s}^{-1}\text{]}$$

- ▶ From before : $v_f = -gT + w \log_e(m_i/m_f)$

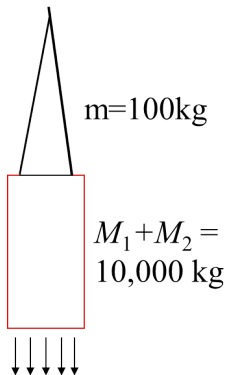
$$m_i = M_1 + M_2 + m ; m_f = 0.25(M_1 + M_2) + m$$

- ▶ Single stage : $v_f = 3.25$ km s^{-1}

- ▶ Earth's escape speed :

$$\frac{1}{2}mv^2 > \frac{GM_em}{R_E} \rightarrow v_{esc} = \sqrt{\frac{2GM_E}{R_E}}$$

$$\rightarrow v_{esc} = 11.2 \text{ km s}^{-1} \text{ (i.e. } v_f < v_{esc}\text{)}$$



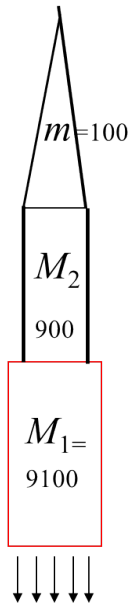
(ii) The 2-stage rocket

Stage 1

- ▶ $|\frac{\Delta m}{\Delta t}| = \alpha \rightarrow T = 0.75M_1/500 \text{ [kg s}^{-1}\text{]}$
- ▶ $v_1 = -gT + w \log_e(m_i/m_f)$
 $m_i = M_1 + M_2 + m ; m_f = 0.25M_1 + M_2 + m$
- ▶ After first stage : $v_f = 2.68 \text{ km s}^{-1}$

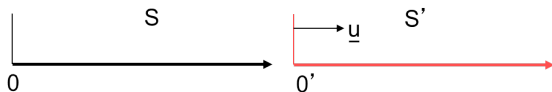
Stage 2

- ▶ $|\frac{\Delta m}{\Delta t}| = \alpha \rightarrow T = 0.75M_2/500 \text{ [kg s}^{-1}\text{]}$
- ▶ $v_2 = v_1 - gT + w \log_e(m_i/m_f)$
 $m_i = M_2 + m ; m_f = 0.25M_2 + m$
- ▶ After second stage :
 $v_2 = 2.68 + 2.80 = 5.48 \text{ km s}^{-1}$
- ▶ Need a 3rd stage, more thrust or less payload to escape



12.3 Non-inertial reference frames

A frame in which Newton's first law is not satisfied - the frame is accelerating (i.e. subject to an external force)

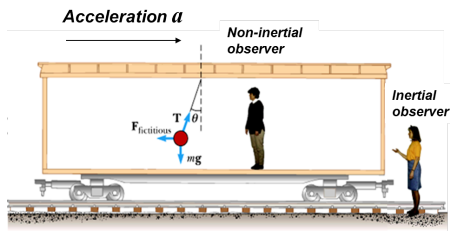


Recall Galilean transformation but now \underline{u} varies with time :

	Galilean transformation	Accelerating frame
Position	$\underline{\mathbf{r}}' = \underline{\mathbf{r}} - \underline{\mathbf{u}}t$	$\underline{\mathbf{r}}' = \underline{\mathbf{r}} - \int \underline{\mathbf{u}}(t) dt$
Velocity	$\underline{\mathbf{v}}' = \underline{\mathbf{v}} - \underline{\mathbf{u}}$	$\underline{\mathbf{v}}' = \underline{\mathbf{v}} - \underline{\mathbf{u}}(t)$
Acceleration	$\frac{d\underline{\mathbf{v}}'}{dt} = \frac{d\underline{\mathbf{v}}}{dt}$	$\frac{d\underline{\mathbf{v}}'}{dt} = \frac{d\underline{\mathbf{v}}}{dt} - \frac{d\underline{\mathbf{u}}}{dt}$
Force on mass	$\underline{\mathbf{F}}'(\underline{\mathbf{r}}') = \underline{\mathbf{F}}(\underline{\mathbf{r}}) = m \frac{d\underline{\mathbf{v}}}{dt}$	$\underline{\mathbf{F}}'(\underline{\mathbf{r}}') = \underline{\mathbf{F}}(\underline{\mathbf{r}}) - m \frac{d\underline{\mathbf{u}}}{dt}$

- ▶ So even if $\underline{\mathbf{F}}(\underline{\mathbf{r}}) = 0$, from NII, there is an *apparent (or "fictitious") force* acting in S' of $\underline{\mathbf{F}}'(\underline{\mathbf{r}}') = -m \frac{d\underline{\mathbf{u}}}{dt}$

12.3.1 Commonplace examples



Accelerating train

- ▶ In the inertial frame

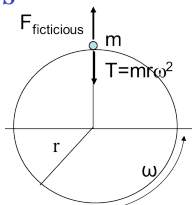
$$\sum F_x = T \sin \theta = ma$$

$$\sum F_y = T \cos \theta - mg = 0$$
- ▶ In the non-inertial frame

$$\sum F'_x = T \sin \theta - F_{\text{fictitious}}$$

$$\sum F'_y = T \cos \theta - mg = 0$$

In NIF need to introduce $F_{\text{fictitious}} = ma$ to explain the displacement of the bob.



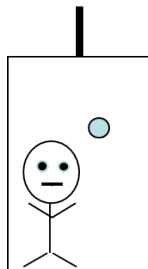
Mass rotating in a circle

- ▶ In the inertial frame
Centripetal acceleration provided by tension in the string
$$T = mr\omega^2$$
- ▶ In the non-inertial frame
The block is at rest and its acceleration is zero

In NIF need $F_{\text{fictitious}} = mr\omega^2$ (centrifugal force) to balance the tension in the string.

12.3.2 Example : accelerating lift

- ▶ First consider the lift in free-fall
- ▶ The ball is “weightless” (stationary or moves at constant velocity) according to an observer in the lift \rightarrow lift becomes an inertial frame (like in deep space) : NI.
- ▶ To this observer the fictitious acceleration $\left(-\frac{d\mathbf{u}}{dt}\right)$ balances the gravitational acceleration



**Free
falling lift**

Observer in accelerating lift

- ▶ Lift plus passenger (total mass M) is now accelerated upwards with force F . Passenger drops ball mass m' from height h ($M \gg m'$).

- ▶ Total force on lift $F_{tot} = F - Mg = Ma$

Acceleration of lift $a = \frac{F}{M} - g$

- ▶ According to passenger in lift frame, *downwards* acceleration of ball is

$$= \left(\frac{F}{M} - g\right) + g = \frac{F}{M} \text{ (downwards)}$$

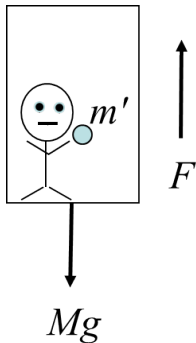
- ▶ Hence weight of ball appears to passenger to be $(F/M) \times m'$

- ▶ Check this:

If $F = 0$, free-fall, ball is weightless

If $F = Mg$, lift is stationary, ball has weight $m'g$

- ▶ Time for ball to reach floor, use $h = \frac{1}{2}at^2 \rightarrow t = \sqrt{\frac{2hM}{F}}$



Observer watching from a frame outside the lift

- ▶ Observer watches ball fall with acceleration g and the lift rise with acceleration $\frac{F}{M} - g$

- ▶ Equate times when ball reaches floor:

$$t = \underbrace{\sqrt{\frac{2h'}{g}}}_{\text{ball falling}} = \underbrace{\sqrt{\frac{2(h-h')}{\frac{F}{M} - g}}}_{\text{lift rising}}$$

- ▶ Solve for h' $\rightarrow h' = \frac{Mgh}{F}$

- ▶ Substitute back

$$\rightarrow t = \sqrt{\frac{2Mh}{F}} \text{ as before.}$$

