CP2 ELECTROMAGNETISM

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LECTURE 21:

ELECTROMAGNETIC WAVES & ENERGY FLOW



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$$\nabla \cdot \mathbf{E} = \frac{\rho}{\varepsilon_0}$$
$$\nabla \cdot \mathbf{B} = 0$$
$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$
$$\frac{1}{\mu_0} \nabla \times \mathbf{B} = \mathbf{J} + \varepsilon_0 \frac{\partial \mathbf{E}}{\partial t}$$

¹ With thanks to Prof Laura Herz

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OUTLINE : 21. ELECTROMAGNETIC WAVES & ENERGY FLOW

- 21.1 Divergence, time derivative, and curl of \underline{E} and \underline{B}
- 21.2 Electromagnetic waves : speed of propagation
- 21.3 Relationship between $\underline{\mathbf{E}}$ and $\underline{\mathbf{B}}$
- 21.4 Electromagnetic wave travelling along the z direction
- 21.5 Characteristic impedance of free space
- 21.6 Polarisation
- 21.7 Energy flow and the Poynting Vector
- 21.8 Example : Poynting Vector for a long resistive cylinder

21.1 Divergence, time derivative, and curl of $\underline{\mathbf{E}}$ and $\underline{\mathbf{B}}$

- ► The divergence of $\underline{\mathbf{E}}$: $\underline{\nabla} \cdot \underline{\mathbf{E}} = \underline{\nabla} \cdot \underline{\mathbf{E}}_{0} \exp\left[i(\omega t \underline{\mathbf{k}} \cdot \underline{\mathbf{r}})\right]$ = $\left[\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z}\right] \cdot \underline{\mathbf{E}}_{0} \exp\left(i(\omega t - k_{x}x - k_{y}y - k_{z}z)\right)$ = $\left[(-i)k_{x}E_{x} + (-i)k_{y}E_{y} + (-i)k_{z}E_{z}\right] \exp\left(i(\omega t - \underline{\mathbf{k}} \cdot \underline{\mathbf{r}})\right)$ = $(-i)\underline{\mathbf{k}} \cdot \underline{\mathbf{E}}$: hence $\underline{\nabla} \equiv -i\underline{\mathbf{k}}$
- The time derivative of $\underline{\mathbf{E}}$: $\frac{\partial}{\partial t} \underline{\mathbf{E}} = \frac{\partial}{\partial t} \underline{\mathbf{E}}_0 \exp \left[i(\omega t \underline{\mathbf{k}} \cdot \underline{\mathbf{r}})\right]$
 - $= i \omega \underline{\mathbf{E}}$: hence $\frac{\partial}{\partial t} \equiv i \omega$
- The curl of $\underline{\mathbf{E}}$:

$$\underline{\nabla} \times \underline{\mathbf{E}} = \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ E_x & E_y & E_z \end{vmatrix} = \begin{pmatrix} \frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} \\ \frac{\partial E_x}{\partial z} - \frac{\partial E_z}{\partial x} \\ \frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} \end{pmatrix} = (-i) \begin{pmatrix} k_y E_z - k_z E_y \\ k_z E_x - k_x E_z \\ k_x E_y - k_y E_x \end{pmatrix} = (-i) \mathbf{\underline{k}} \times \mathbf{\underline{E}} \quad \& \text{ again } \mathbf{\underline{\nabla}} \equiv -i \mathbf{\underline{k}}$$

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21.2 Electromagnetic waves : speed of propagation

► To get speed of propagation, substitute $\underline{\mathbf{E}} = \underline{\mathbf{E}}_0 \exp(i(\omega t - \underline{\mathbf{k}} \cdot \underline{\mathbf{r}}))$ into the wave equation $\underline{\nabla}^2 \underline{\mathbf{E}} = \epsilon_0 \mu_0 \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}$

► Use
$$\underline{\nabla} \equiv -i\underline{\mathbf{k}} \rightarrow \underline{\nabla}^2 \equiv (-i\underline{\mathbf{k}})^2 = -k^2$$

 $\frac{\partial}{\partial t} \equiv i\omega \rightarrow \frac{\partial^2}{\partial t^2} \equiv (i\omega)^2 = -\omega^2$
► $-k^2 \underline{\mathbf{E}}_0 \exp(i(\omega t - \underline{\mathbf{k}} \cdot \underline{\mathbf{r}})) = -\omega^2 \epsilon_0 \mu_0 \underline{\mathbf{E}}_0 \exp(i(\omega t - \underline{\mathbf{k}} \cdot \underline{\mathbf{r}}))$
 $\rightarrow k^2 = \omega^2 \epsilon_0 \mu_0$

Fields of this form are solutions to the wave equation with velocity of propagation :

$$c = \frac{\omega}{k} = \frac{1}{\sqrt{\epsilon_0 \, \mu_0}} = 3 \times 10^8 \, \mathrm{m \, s^{-1}}$$

i.e. the speed of light $\rightarrow\,$ speed of an EM wave in vacuum

21.3 Relationship between $\underline{\mathbf{E}}$ and $\underline{\mathbf{B}}$

• Substitute $\underline{\mathbf{E}} = \underline{\mathbf{E}}_0 \exp(i(\omega t - \underline{\mathbf{k}} \cdot \underline{\mathbf{r}}))$ into Maxwell eqn's :

$$\underline{\nabla} \cdot \underline{\mathbf{E}} = -i\,\underline{\mathbf{k}} \cdot \underline{\mathbf{E}} = \mathbf{0}$$

$$\underline{\nabla} \cdot \underline{\mathbf{B}} = -i \, \underline{\mathbf{k}} \cdot \underline{\mathbf{B}} = \mathbf{0}$$

Hence $\underline{\mathbf{k}} \cdot \underline{\mathbf{E}} = \mathbf{0}$ and $\underline{\mathbf{k}} \cdot \underline{\mathbf{B}} = \mathbf{0}$

- ► Electric and magnetic fields in vacuum are *perpendicular* to direction of propogation → *EM waves are transverse*
- ► Substitute into Faraday's Law : $\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$ $-i\mathbf{k} \times \mathbf{E} = -i\omega \mathbf{B} \rightarrow \mathbf{B} = \frac{1}{2}\mathbf{k} \times \mathbf{E}$

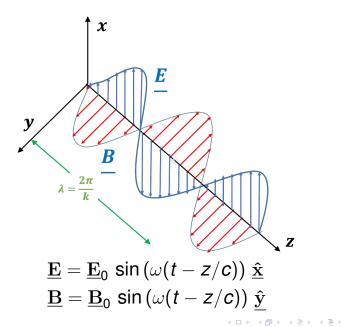
Substitute into Ampere's Law :
$$\underline{\nabla} \times \underline{\mathbf{B}} = \mu_0 \epsilon_0 \frac{\partial \underline{\mathbf{E}}}{\partial t}$$

$$-i\underline{\mathbf{k}} \times \underline{\mathbf{B}} = i\,\omega\,\mu_0\,\epsilon_0\,\underline{\mathbf{E}} \quad \rightarrow \qquad \underline{\mathbf{E}} = -\frac{c^2}{\omega}\,\underline{\mathbf{k}} \times \underline{\mathbf{B}}$$

- $\underline{\mathbf{E}}, \underline{\mathbf{B}} \& \underline{\mathbf{k}}$ are mutually orthogonal (NB. $\underline{\mathbf{k}} \times \underline{\mathbf{B}} = kB \sin \frac{\pi}{2} \underline{\hat{\mathbf{E}}}$)
- $\underline{\mathbf{E}}$ and $\underline{\mathbf{B}}$ are in phase and lie in the plane of the wavefront
- Field magnitude ratio :

$$|\underline{\mathbf{E}}|/|\underline{\mathbf{B}}| = \frac{c^2}{\omega} \mathbf{k} = \mathbf{C} = \frac{1}{\sqrt{\mu_0 \epsilon_0}}$$

21.4 Electromagnetic wave travelling along the z direction



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21.5 Characteristic impedance of free space

- Take the ratio $Z = \frac{|\mathbf{E}|}{|\mathbf{H}|}$ where $|\mathbf{H}| = \frac{1}{\mu_0} |\mathbf{B}|$
- *Z* has units $[V m^{-1}] / [A m^{-1}] = Ohms.$
- Z is called the characteristic impedance of free space

$$Z = \mu_0 \, \frac{|\mathbf{E}|}{|\mathbf{B}|} = \mu_0 \, \mathbf{c} = \frac{\mu_0}{\sqrt{\mu_0 \, \epsilon_0}} = \sqrt{\frac{\mu_0}{\epsilon_0}} = 376.7 \, \Omega$$

Electromagnetic waves : summary

In vacuum, free of charge or currents (
$$\rho$$
, J = 0):
 $\nabla \cdot \mathbf{E} = 0$
 $\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$
 $\nabla \cdot \mathbf{B} = 0$
 $\nabla \times \mathbf{B} = \varepsilon_0 \mu_0 \frac{\partial \mathbf{E}}{\partial t}$
 $\nabla^2 \mathbf{E} = \varepsilon_0 \mu_0 \ddot{\mathbf{B}}$
 $\nabla^2 \mathbf{B} = \varepsilon_0 \mu_0 \ddot{\mathbf{B}}$
Wave equations in E, B!
Electromagnetic waves propagate in free space:
Plane EM wave fronts: $\mathbf{E} = \mathbf{E}_0 \exp\{i(\omega t - \mathbf{k} \cdot \mathbf{r})\}$ with wavelength $\lambda = \frac{2\pi}{k}$
Propagation velocity of wave fronts: $c = \frac{\omega}{k} = \frac{1}{\sqrt{\varepsilon_0 \mu_0}} = 3 \times 10^8 \,\mathrm{m \, s^{-1}}$
Relationship between E and B:
(in phase and mutually orthogonal $\mathbf{B} = \frac{\mathbf{k}}{\omega} \times \mathbf{E}$
 $\mathbf{E} = -c^2 \frac{\mathbf{k}}{\omega} \times \mathbf{B}$
 $|\mathbf{E}| = c$
with wave vector k)
Impedance of free space: $Z = \frac{|\mathbf{E}|}{|\mathbf{B}|/\mu_0} = \sqrt{\frac{\mu_0}{\varepsilon_0}} = 376.7 \,\Omega$

21.6 Polarisation

- Linearly (or plane) polarised wave :
 E has one specific orientation
- Circularly polarised wave :

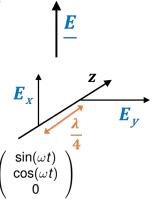
Two linear components of *E* superimposed at a right angle and phase shifted by $\pi/2$

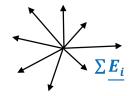
$$\begin{pmatrix} E \\ 0 \\ 0 \end{pmatrix} \sin(\omega t) + \begin{pmatrix} 0 \\ E \\ 0 \end{pmatrix} \sin(\omega t + \pi/2) = E \begin{pmatrix} 0 \\ E \\ 0 \end{pmatrix}$$
Elliptically polarised wave :

As above but with unequal amplitudes

Unpolarised :

E superimposed with all orientations (with no fixed phase relationships between components)





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21.7 Energy flow and the Poynting Vector

 Recall : Energy of the electric field U_e = ∫_ν ¹/₂ ε₀ E² dν Energy of the magnetic field U_m = ∫_ν ¹/_{2μ₀} B² dν
 Total EM energy in volume ν :

$$U = \int_{\nu} \underbrace{\frac{1}{2} \left(\epsilon_{0} \underline{\mathbf{E}} \cdot \underline{\mathbf{E}} + \frac{1}{\mu_{0}} \underline{\mathbf{B}} \cdot \underline{\mathbf{B}} \right)}_{\text{Energy density}} d\nu$$

$$= \ln \text{ free space } (\underline{\mathbf{J}} = 0, \rho = 0) \\ \underbrace{\frac{\partial \underline{\mathbf{E}}}{\partial t} = \frac{1}{\mu_{0}\epsilon_{0}} \underline{\nabla} \times \underline{\mathbf{B}}}_{\text{Ampere's Law}} ; \underbrace{\frac{\partial \underline{\mathbf{B}}}{\partial t} = -\underline{\nabla} \times \underline{\mathbf{E}}}_{\text{Faraday's Law}}$$

$$= \operatorname{Calculate \ the \ rate \ of \ change \ of \ energy \ in \ \nu} \\ \frac{d\underline{U}}{dt} = \int_{\nu} \left(\epsilon_{0} \underline{\mathbf{E}} \cdot \underline{\dot{\mathbf{E}}} + \frac{1}{\mu_{0}} \underline{\mathbf{B}} \cdot \underline{\dot{\mathbf{B}}} \right) d\nu$$

Energy flow and the Poynting Vector continued

- ► Energy flow out of volume \mathcal{V} per unit time : $\frac{dU}{dt} = -\frac{1}{\mu_0} \int_{\mathcal{V}} \nabla \cdot (\mathbf{\underline{E}} \times \mathbf{\underline{B}}) \, d\mathcal{V} \qquad \mathbf{S}$
- Apply the divergence theorem :

$$\frac{dU}{dt} = -\oint_{S} \underbrace{(\frac{1}{\mu_{0}} \mathbf{\underline{E}} \times \mathbf{\underline{B}})}_{\text{Pownting Vector } \mathbf{N}} \cdot d\mathbf{\underline{a}}$$

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$$\frac{dU}{dt} = -\oint_{\mathcal{S}} \mathbf{\underline{N}} \cdot \mathbf{d}\mathbf{\underline{a}} \quad \text{where} \quad \mathbf{\underline{N}} = \frac{1}{\mu_0} \mathbf{\underline{E}} \times \mathbf{\underline{B}}$$

Poynting vector \underline{N} is the power per unit area flowing through the surface bounded by volume \mathcal{V} . (It also gives the direction of flow). Units of \underline{N} : $[W m^{-2}]$

► For EM waves, the intensity is the time-average of $|\mathbf{N}|$ $\Im = < |\mathbf{N}| >= \frac{1}{\mu_0} E_0 B_0 \underbrace{<\cos^2(\omega t - \mathbf{k} \cdot \mathbf{r}) >}_{1/2} = \frac{1}{2\mu_0 c} E_0^2$

21.8 Example : Poynting Vector for a long resistive cylinder

- Calculate Poynting Vector at the surface of the wire with applied potential difference V and current I
 <u>N</u> = ¹/_{µ0}<u>E</u> × <u>B</u>
- Electric field along wire axis : $E = V/\ell$

Magnetic flux density at wire surface : $\oint \mathbf{\underline{B}} \cdot d\underline{\ell} = \mathbf{B} \cdot 2\pi \mathbf{a} = \mu_0 I$

(note that this is *tangential* - along circumference)

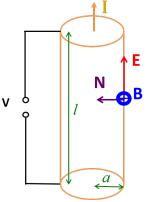
N = ¹/_{µ₀} ^V/_ℓ ^{µ₀I}/_{2πa}
 (in radial direction pointing *inwards* - i.e. wire heats up !)

• Hence
$$N = (VI)/2\pi\ell a$$

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surface area

- Total power dissipated in wire : $P = \int_{S} \mathbf{N} \cdot d\mathbf{a} = V I$
 - as expected from circuit theory.



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Poynting Vector : summary

